

A. Theoretical Fundamentals of Airborne Gravimetry, Parts I and II (Monday, 23 May 2016)

I. Introduction - Airborne Gravity Data Acquisition

II. Elemental Review of Physical Geodesy

III. Basic Theory of Moving-Base Scalar Gravimetry

IV. Overview of Airborne Gravimetry Systems

B. Theoretical Fundamentals of Gravity Gradiometry and Inertial Gravimetry (Thursday, 26 May 2016)

V. Theoretical Fundamentals of Inertial Gravimetry

VI. Theoretical Fundamentals of Airborne Gradiometry

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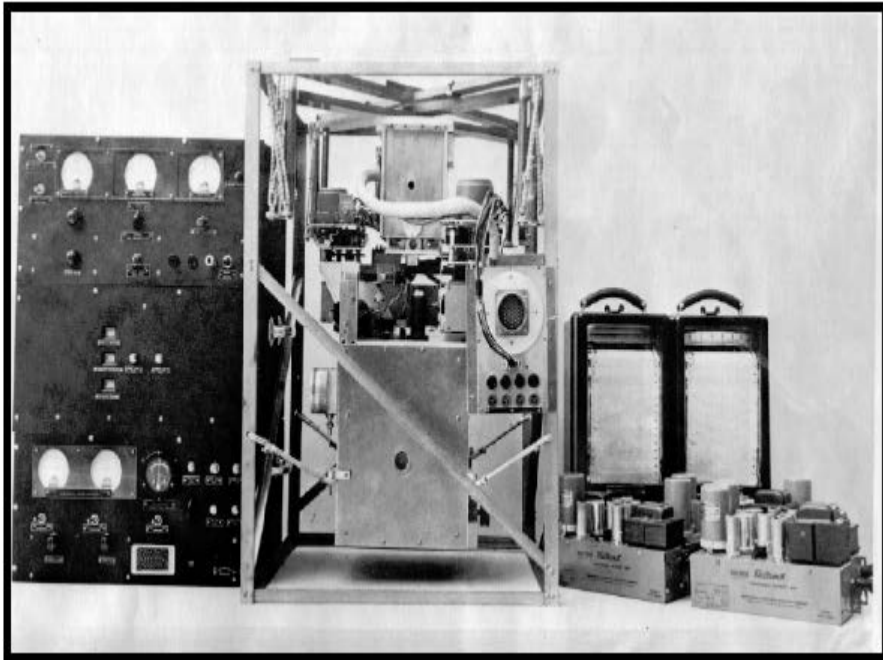
I. Introduction - Airborne Gravity Data Acquisition

- **A very brief history of airborne gravimetry**
- **Why airborne gravimetry?**

A Brief History of Airborne Gravimetry

- **Natural evolution of successes in 1st half of 20th century with **ocean-bottom**, **submarine**, and **shipboard** gravimeters operating in dynamic environments**
- airborne systems promised rapid, if not highly accurate, regional gravity maps for exploration reconnaissance and military geodetic applications
- **Special challenges**
- critical errors are functions of **speed** and **speed-squared**
- difficulty in accurate **altitude & vertical acceleration** determination
- trade accuracy for acquisition speed
- **1958: First fixed-wing airborne gravimetry test (Thompson and LaCoste 1960)**
- 5-10 minute average, 10 mgal accuracy
- high altitude, 6-9 km
- **Further tests by exploration concerns**
- **LaCoste & Romberg, Austin TX**
- **Gravity Meter Exploration Co., Houston, TX**
- **10 mGal accuracy, 3 minute averages (Nettleton et al. 1960)**

First Airborne Gravity Test – Air Force Geophysics Lab 1958



Instruction Manual LaCoste Romberg Model "S" Air-Sea Dynamic Gravity Meter, 2002; with permission

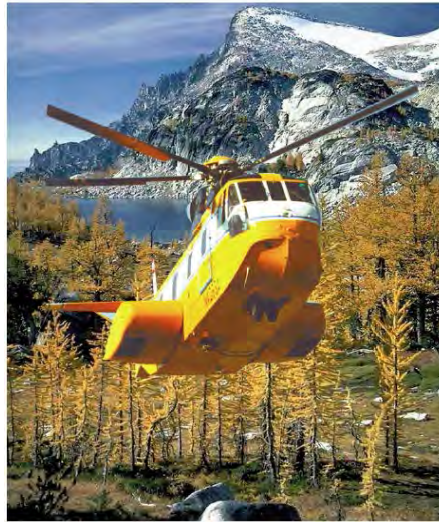
- **The first LaCoste-Romberg Model "S" Air-Sea Gravimeter**

- **Doppler navigation system – elevation above mean sea level determined from the tracking range data**

- **flights over an Askania camera tracking range at Edwards Air Force Base**

- **KC-135 jet tanker**

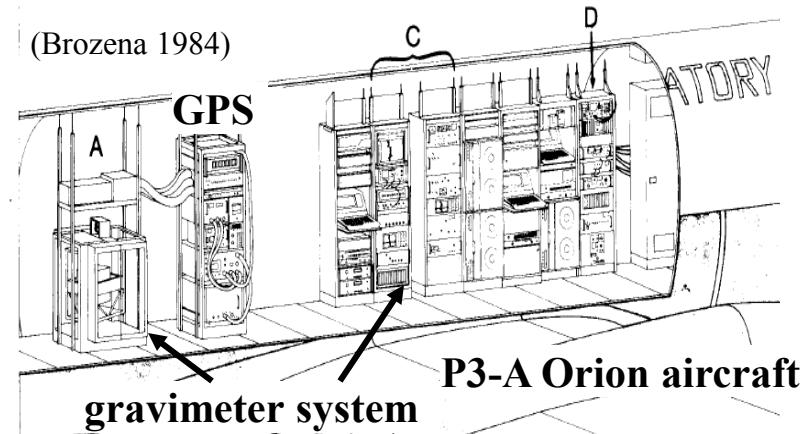
First Successful Helicopter Airborne Gravimetry Test 1965



- **Carson Services, Inc. (Carson Helicopter)**
 - gimbal-suspended LaCoste and Romberg Sea gravimeter
 - 5 mGal accuracy, hovering at 15 m altitude (Gumert 1998)
 - Navy sponsored
- **Further tests and development by exploration companies**
 - principally, Carson Services throughout the 1960s and 1970s

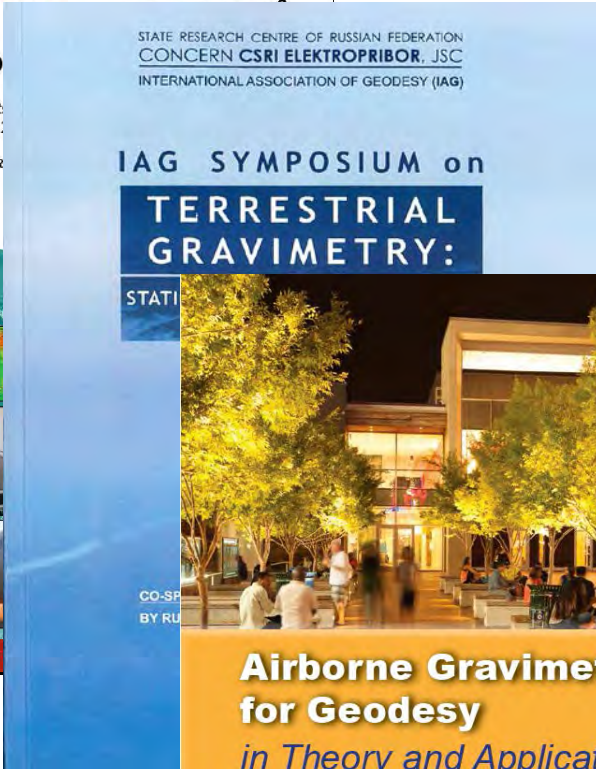
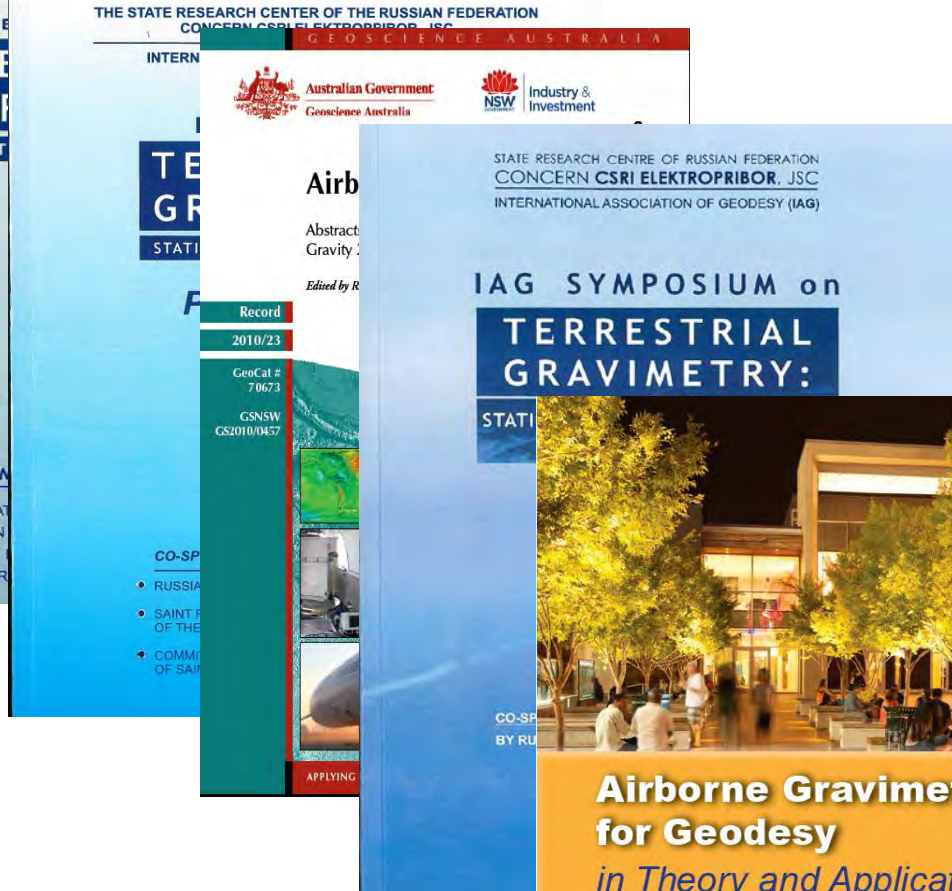
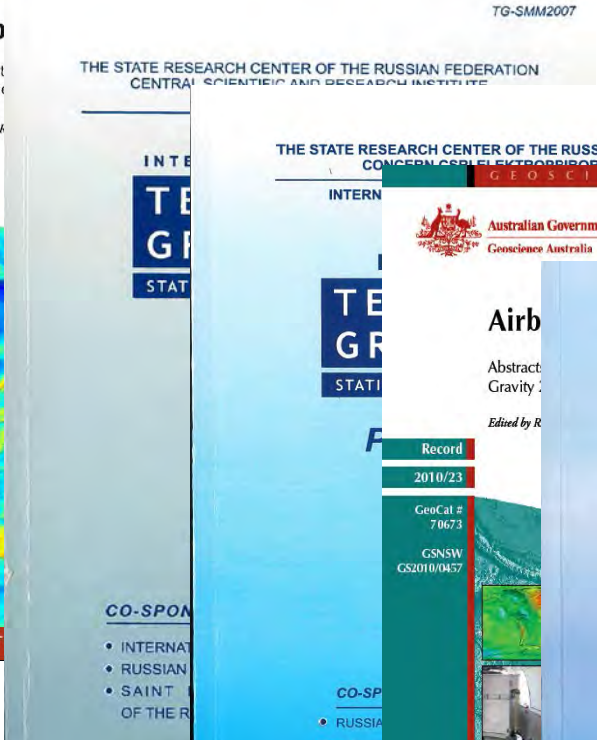
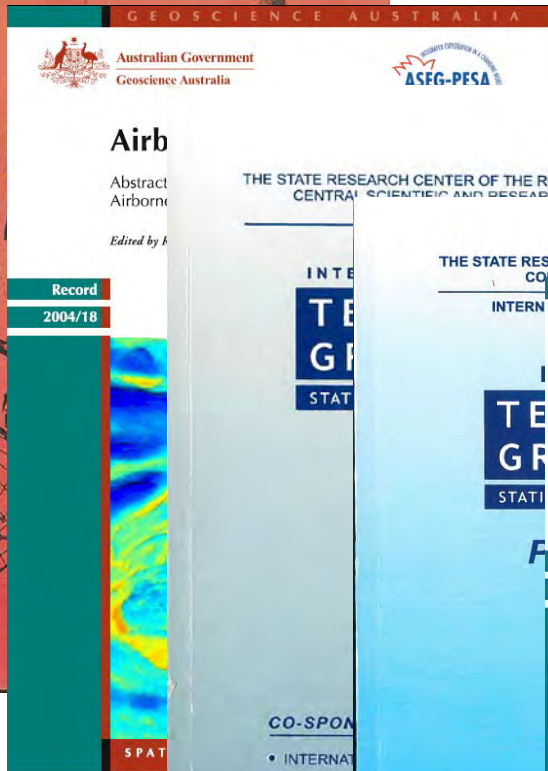
Rapid Development with Advent of GPS (1980s and 1990s)

- **Naval Research Laboratory**
 - John Brozena
- **National Survey and Cadastre of Denmark (DKM) – Rene Forsberg**
- **Academia (in collaboration with industry and government)**
 - University of Calgary (K.P. Schwarz)
 - University FAF Munich (G. Hein)
 - Swiss Federal Institute of Technology (E.E. Klingele)
 - Lamont-Doherty Earth (Geological) Observatory (R. Bell)
 - ...
- **Industry ...**



Twin-Otter Aircraft

Dedicated International Symposia & Workshops

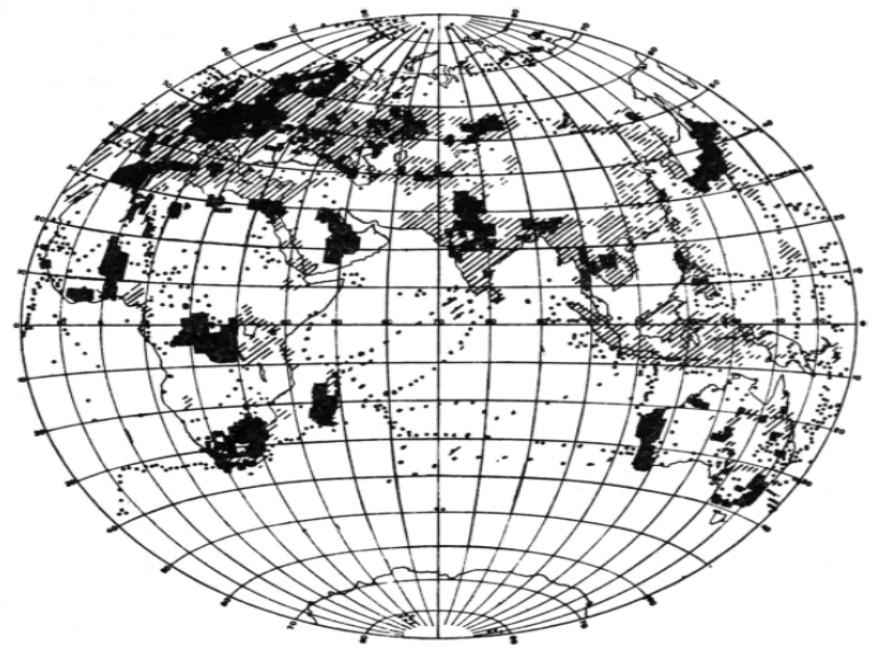


Airborne Gravimetry for Geodesy in Theory and Application Summer School
May 23-27, 2016

Free-Air Gravity Disturbance (mGal)
Canada
Scale: 0 100 200 400 Kilometers

The Need for Global Gravity Data

- Gravity data until the early 1960s were obtained primarily by point measurements on land and along some ship tracks.
 - map of data archive of 1963 (Kaula 1963)



1990s – More Data, Still Many Gaps

- **Greater uniformity, but only at relatively low resolution**
 - map of terrestrial $1^\circ \times 1^\circ$ anomaly archive of 1990 (Rapp and Pavlis 1990)

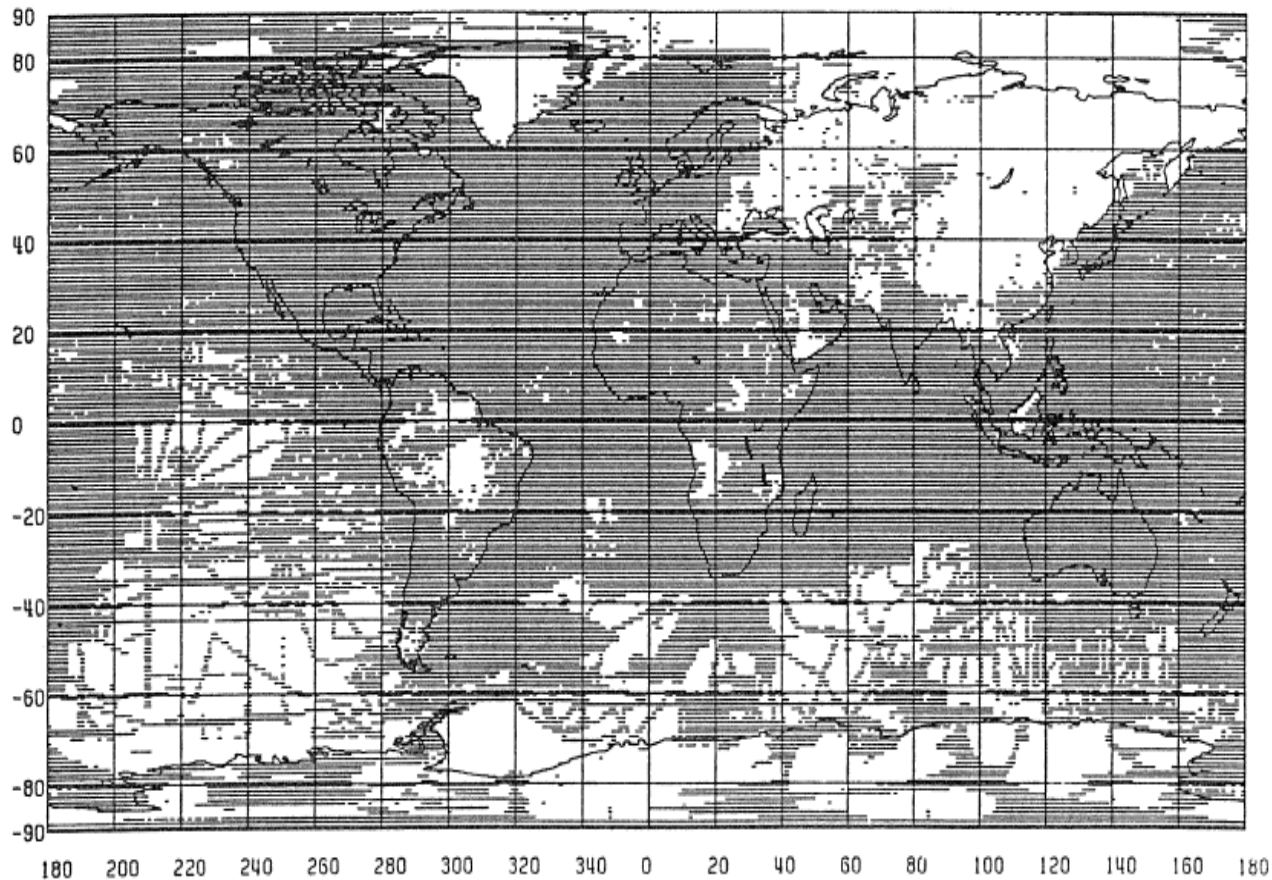
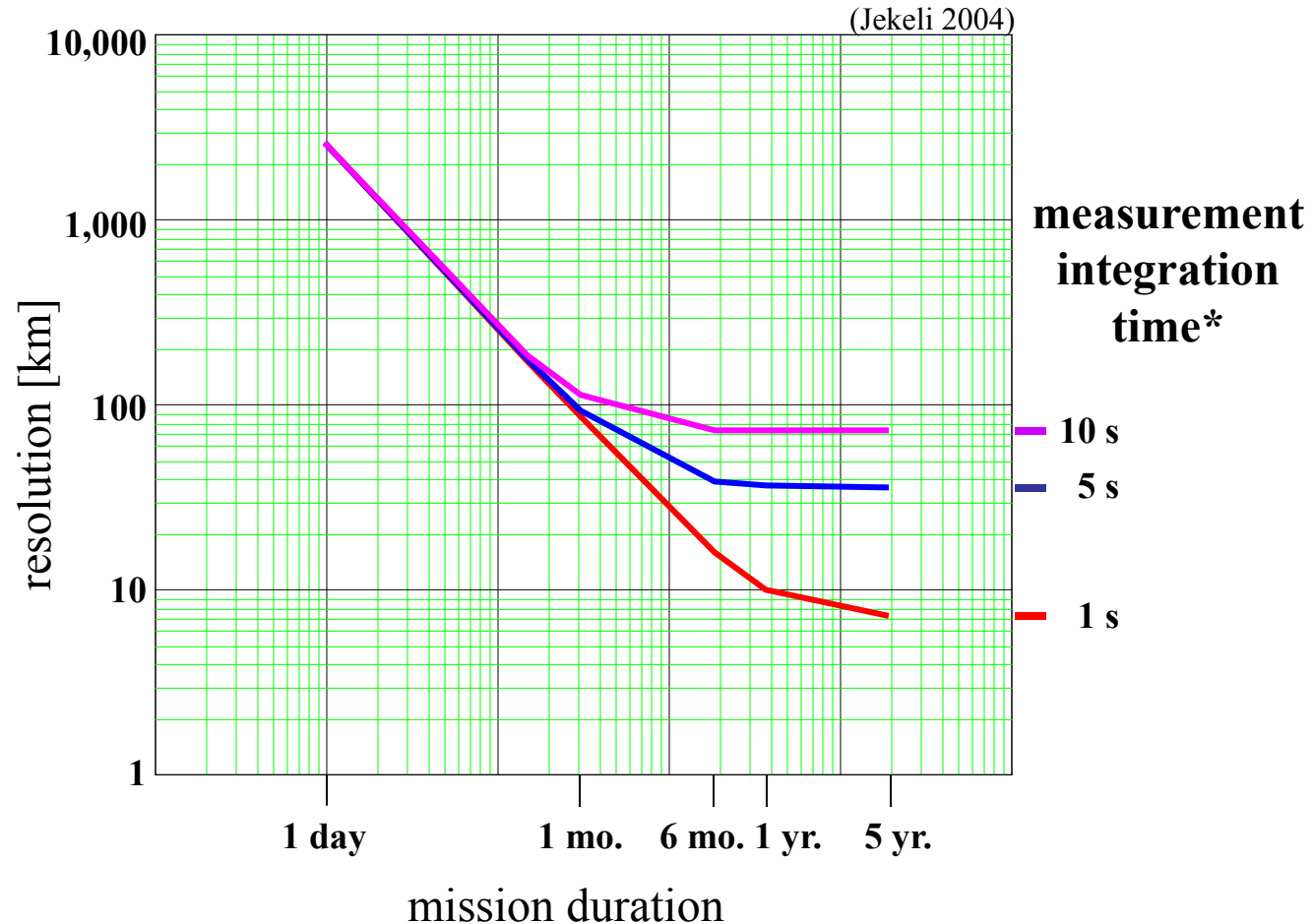


Fig. 3. Location of 45,126 $1^\circ \times 1^\circ$ free-air anomalies in the OSU July 1989 file excluding geophysically predicted values.

Why Airborne Gravimetry?

Satellite Resolution vs Mission Duration and Integration Time

- Satellite-derived gravitational models are **limited in spatial resolution** because of high inherent satellite **speed**
- Only **airborne gravimetry** yields higher resolution efficiently

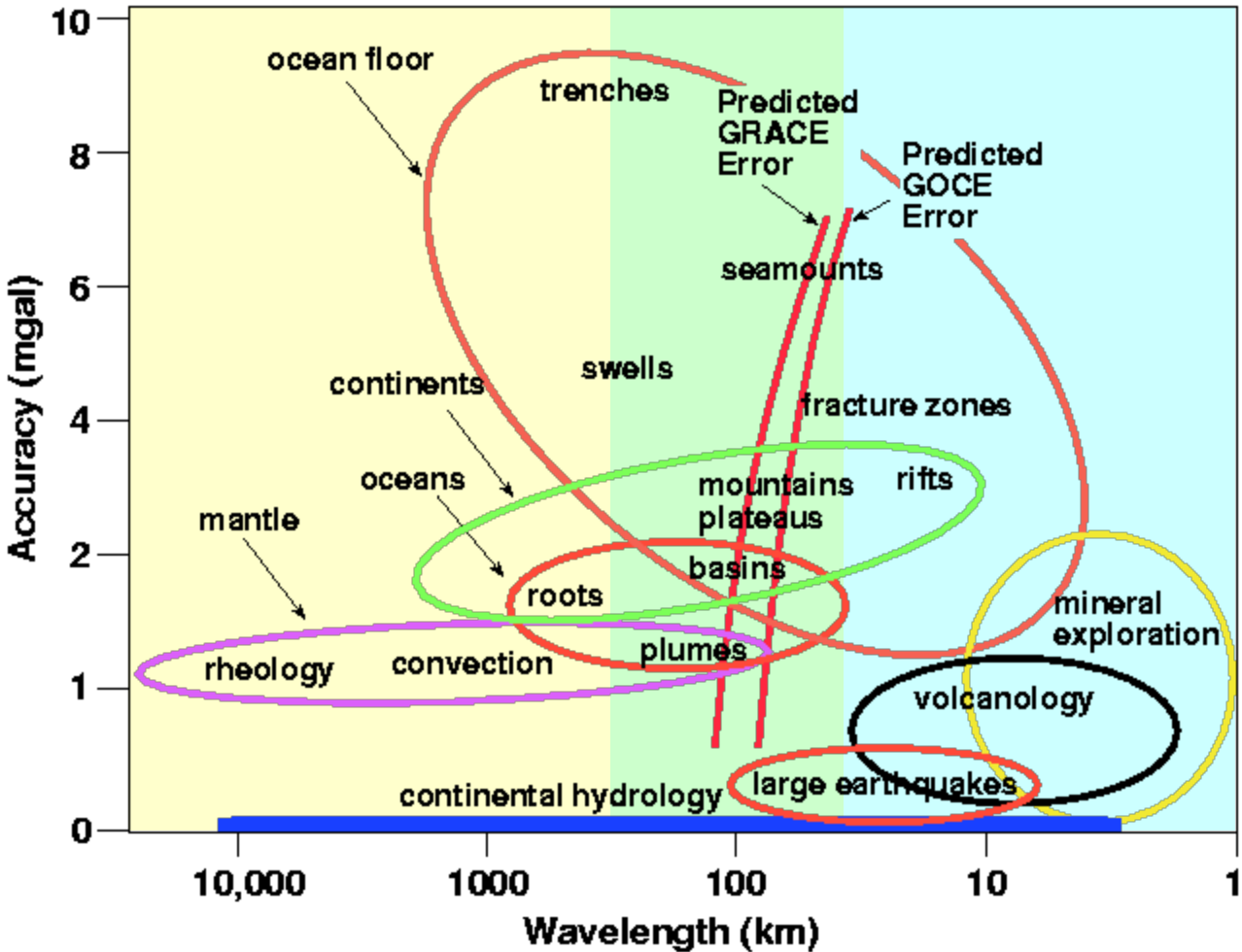


*

CHAMP: 10 s	GOCE: 5 - 10 s
GRACE: 5 s	Laser-interferometry GRACE follow-on: 1 - 10 s

Gravity Resolution vs Accuracy Requirements in Geophysics

Satellite Gravimetry/Gradiometry Airborne Gravimetry/Gradiometry



Geodetic Motivation

sea-level rise



River flooding



coastal flooding from hurricane (Sandy)



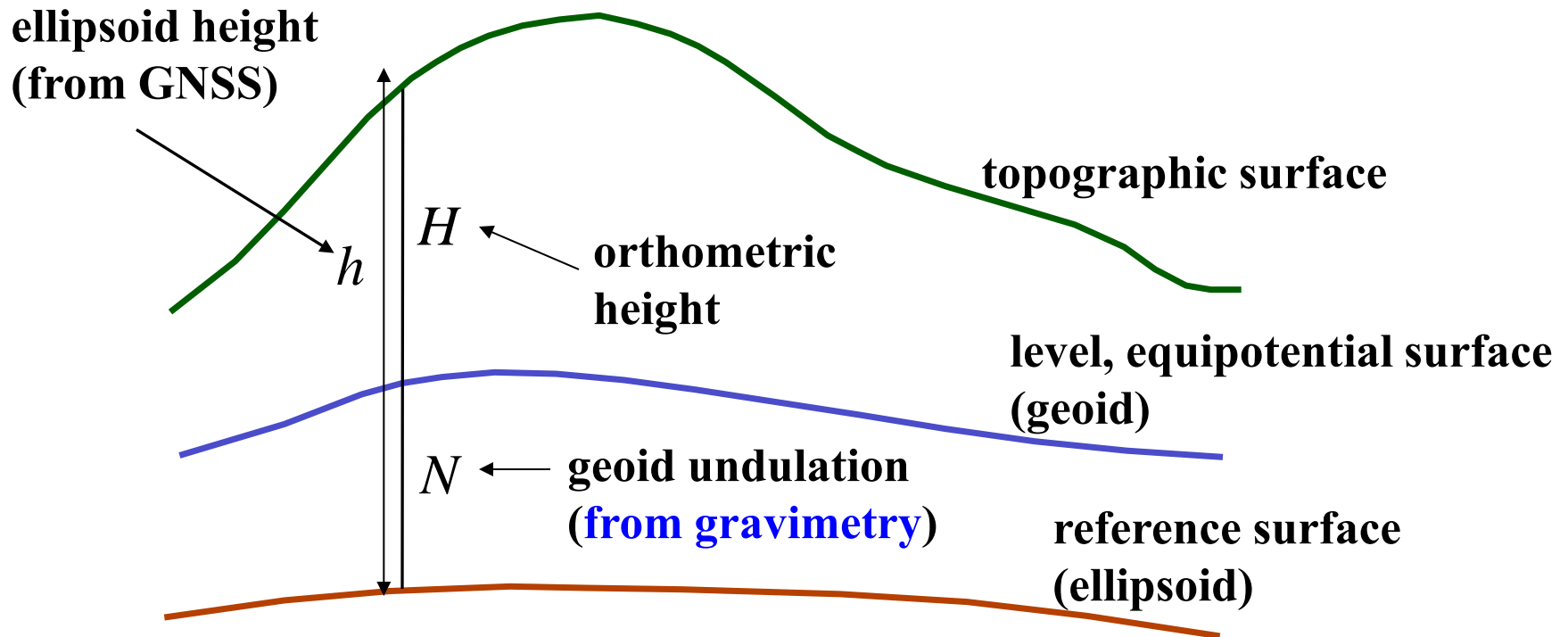
[http://www.fourwinds10.net/resources/uploads/images/missouri%20river%20flooding\(1\).jpg](http://www.fourwinds10.net/resources/uploads/images/missouri%20river%20flooding(1).jpg)

http://theistgringo.com/2012/11/northeast-united-states-future-climate-change-to-the-year-2099/homes-flooded-nj-hurricane-sandy-oct-2012_620_422_s_c1-2/

http://images.nationalgeographic.com/wp/f/media-live/photos/000/166/cache/article-sea-level-rise_16648_600x450.jpg

GNSS and Geopotential

- Traditional height reference surface: **equipotential surface (geoid)**
 - needed for determining and monitoring the **flow of water**, from flood control to sea level rise
 - replace arduous spirit leveling with GNSS: $H = h - N$





II. Elemental Review of Physical Geodesy

- **Gravitational potential, gravity**
- **Normal gravity**
- **Disturbing potential, gravity anomaly, deflection of the vertical**
- **Geoid determination aspects**

Basic Definitions

- **Gravitational potential, V**

- due to mass attraction

- **gravitational** acceleration: $\mathbf{g} = \nabla V$

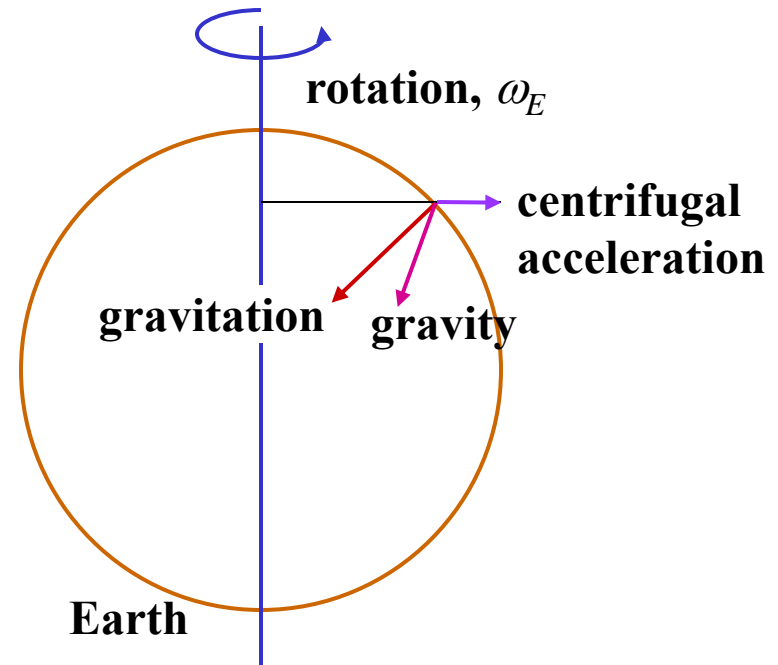
- **Centrifugal “potential”, ϕ**

- due to Earth’s rotation

- centrifugal acceleration: $\mathbf{a}_{\text{cent}} = \nabla \phi$

- **Gravity potential, $W = V + \phi$**

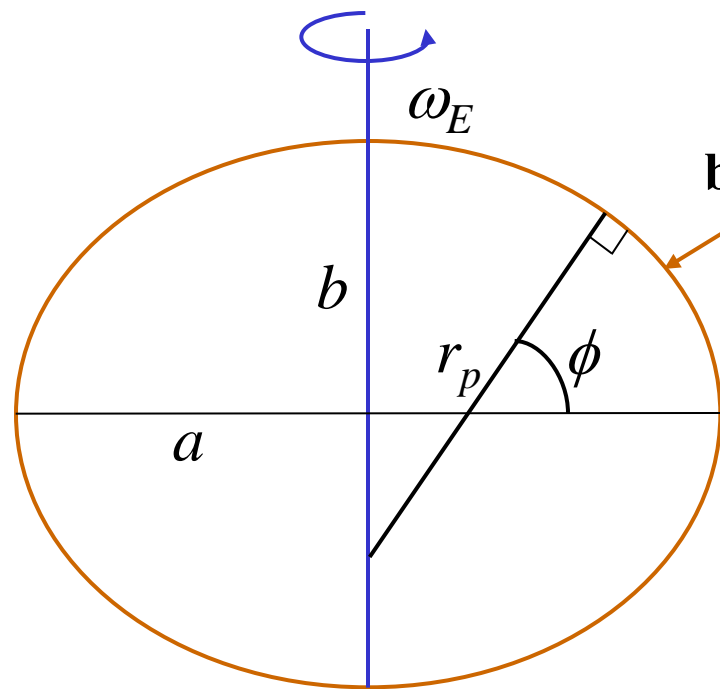
- **gravity** acceleration: $\bar{\mathbf{g}} = \mathbf{g} + \mathbf{a}_{\text{cent}}$



Physical geodesy makes the distinction between **gravitation** and **gravity**, especially in terrestrial gravimetry

Normal Gravitational Potential

- **Mathematically simple potential and boundary**
 - approximates Earth's potential and geoid to about 5 ppm
 - approximates Earth's gravity to about 50 ppm
 - rotates with the Earth



boundary = rotational ellipsoid

centrifugal potential on ellipsoid

$$V|_{\text{boundary}} = U_0 - \underbrace{\frac{1}{2} \omega_e^2 r_p^2 \cos^2 \phi}$$

- **boundary function is chosen so that gravity potential on boundary is a constant, U_0**

Normal Gravity Potential

- Expressed as spherical harmonic series in spherical coordinates

$$U(r, \theta) = V(r, \theta) + \phi(r, \theta)$$
$$= \frac{GM}{a} \sum_{n=0}^{\infty} \left(\frac{a}{r}\right)^{2n+1} C_{2n}^N P_{2n}(\cos \theta) + \frac{1}{2} \omega_e^2 r^2 \sin^2 \theta$$

– closed expression exists in ellipsoidal coordinates

– C_{2n}^N depends on only 4 parameters: ω_e, C_2^N, a, GM

– (e.g., WGS84 parameters)

- Normal gravity vector: $\gamma = \nabla U$

$$\omega_e = 7.292115 \times 10^{-5} \text{ rad/s}$$

$$C_2^N = -0.484166774985 \times 10^{-3}$$

$$a = 6378137. \text{ m}$$

$$GM = 3.986004418 \times 10^{14} \text{ m}^3/\text{s}^2$$

Gravity Disturbance and Anomaly

• **Disturbing potential:** $T = W - U$ ($W = \text{total gravity potential}$)

• **Gravity disturbance vector:** $\delta \mathbf{g} = \nabla W - \nabla U = \mathbf{g} - \boldsymbol{\gamma}$

– **gravity disturbance:** $\delta g = |\mathbf{g}| - |\boldsymbol{\gamma}|$

– **in n -frame (North-East-Down):** $\delta \mathbf{g}^n = \begin{pmatrix} g_N - \gamma_N \\ g_E \\ g_D - \gamma_D \end{pmatrix} \approx \begin{pmatrix} g_N \\ g_E \\ g_D - \gamma_D \end{pmatrix}$

◦ due to symmetry, $\gamma_E = 0$

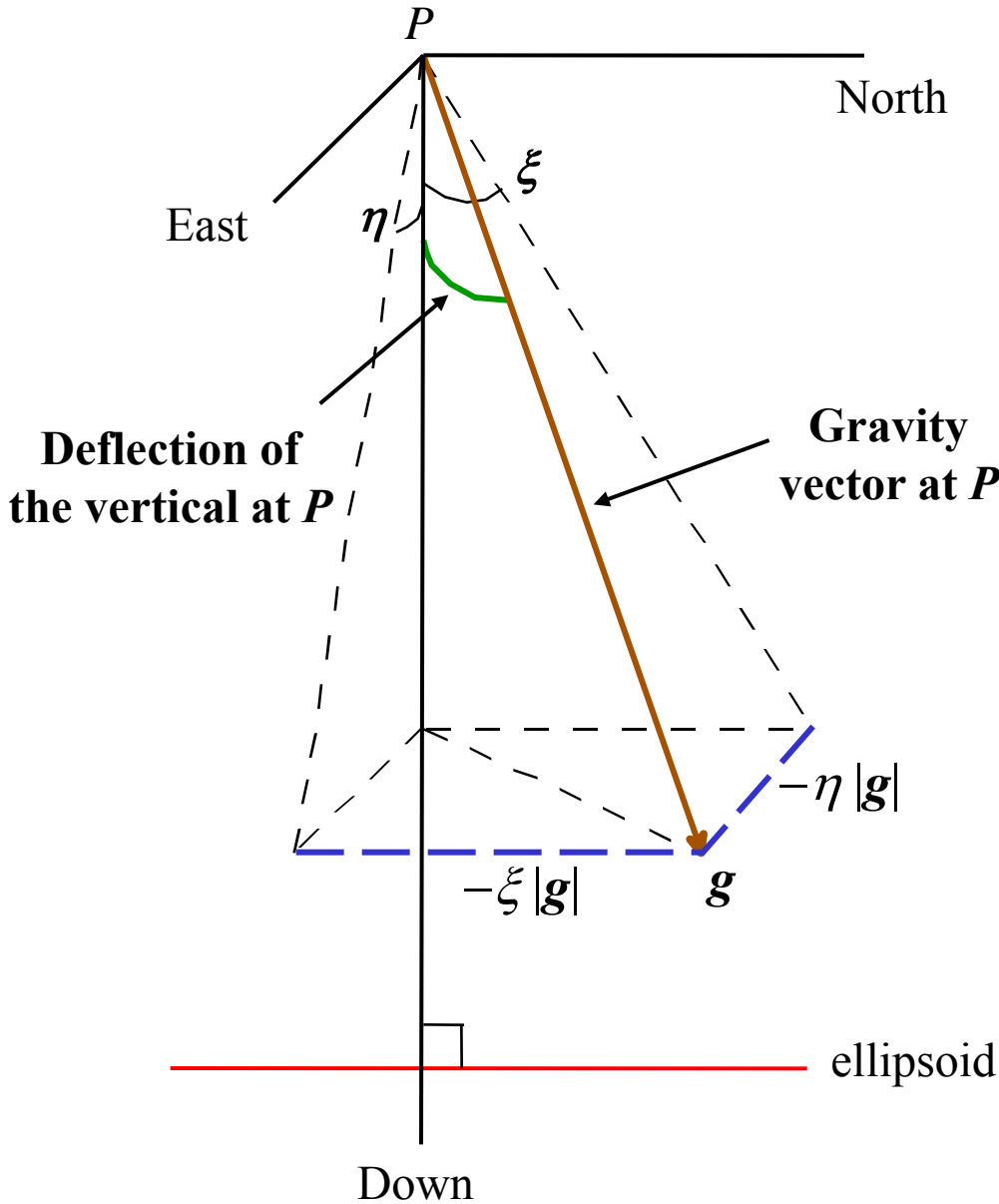
◦ near Earth's surface, $\gamma_N \approx 0$

• **Gravity anomaly vector:** $\Delta \mathbf{g}_P = \nabla W_P - \nabla U_Q = \mathbf{g}_P - \boldsymbol{\gamma}_Q$

– P and Q are points on the ellipsoid normal such that $W_P = U_Q$

– **gravity anomaly:** $\Delta g_P = |\mathbf{g}_P| - |\boldsymbol{\gamma}_Q|$

Deflection of the Vertical



$\xi =$ north deflection

$\eta =$ east deflection

$$\nabla T = \delta \mathbf{g}^n \approx \begin{pmatrix} -\xi g \\ -\eta g \\ \delta g \end{pmatrix} \approx \begin{pmatrix} g_N \\ g_E \\ g_D - \gamma_D \end{pmatrix}$$

- linear approximation
- signs agree with convention of **astronomic** deflection of the vertical

Geoid Determination

• **Brun's Formula:** $N_{P_0} = \frac{1}{\gamma_{Q_0}} T_{P_0} + N_0$ where N_0 is a height datum offset

• **Boundary-value Problem:** $\nabla^2 T = 0$ above geoid (by assumption)

$$\Delta g_{P'_0} = -\frac{\partial T}{\partial h} \Big|_{P'_0} + \frac{1}{\gamma_{Q_0}} \frac{\partial \gamma}{\partial h} \Big|_{Q_0} T_{P'_0} \quad \text{boundary condition}$$

$$\Delta g_{P'}, \Delta g_{P'_a}$$

+

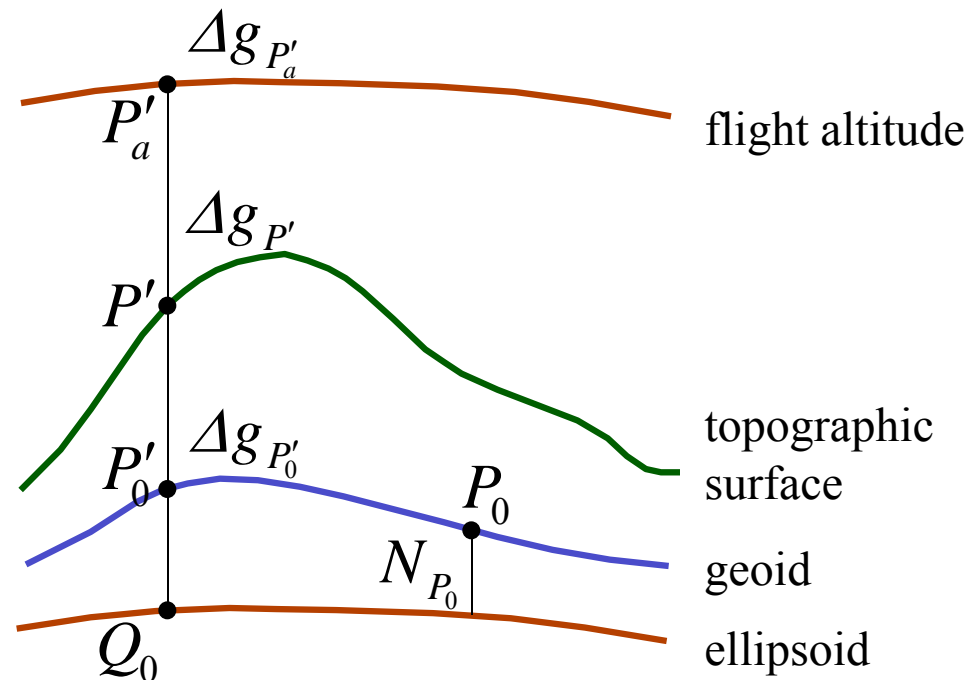
gravity reductions

||

$$\Delta g_{P'_0}$$

• **Stokes's formula**

$$N_{P_0} = N_0 + \frac{R}{4\pi\gamma_{Q_0}} \iint_{\Omega} \Delta g_{P'_0} S(\psi_{P_0, P'_0}) d\Omega$$



Details

- **Gravity reductions to satisfy the boundary-value conditions**
 - re-distribution of topographic mass; consequent indirect effect
 - downward continuation (various methods)
- **Ellipsoidal corrections**
 - account for spherical approximation of geoid, boundary condition
- **Include existing spherical harmonic model (satellite-derived)**
 - remove-compute-restore techniques
- **Back to Motivation**
 - use airborne gravimetry to **improve spatial resolution of data** (boundary values) – few km to 200 km wavelengths



III. Basic Theory of Moving-Base Scalar Gravimetry

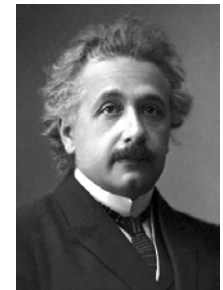
- **Fundamental laws of physics and the gravimetry equation**
- **Coordinate frames**
- **Mechanizations and methods of scalar gravimetry**
- **Rudimentary error analyses**

Fundamental Physical Laws

- **Moving-Base Gravimetry and Gradiometry are based on 3 fundamental laws in physics**
 - **Newton's Second Law of Motion**
 - **Newton's Law of Gravitation**
 - **Einstein's Equivalence Principle**
- **Laws are expressed in an inertial frame**
- **General Relativistic effects are not yet needed**
 - however, the interpretation of space in the theory of general relativity is used to distinguish between applied and gravitational forces

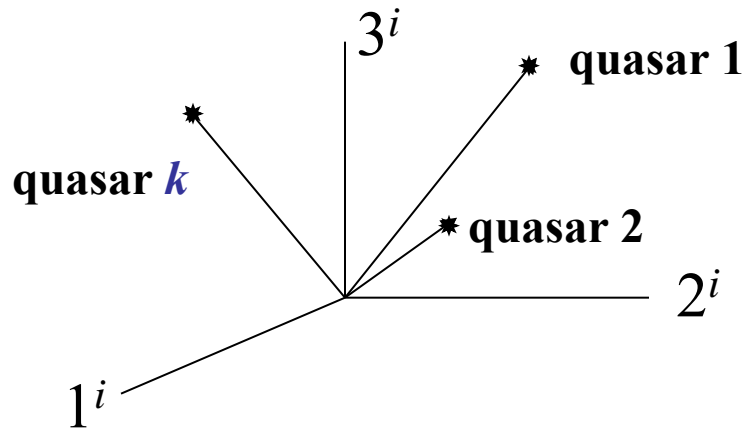


Issac Newton
1643 - 1727



Albert Einstein
1879 - 1955

Inertial Frame



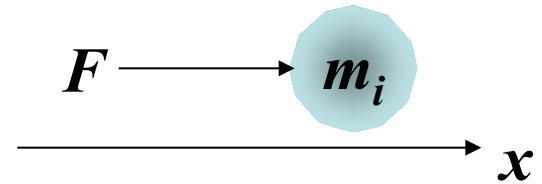
notation convention:
- axis identified by number
- superscript identifies frame

- The realization of a **system** of coordinates that **does not rotate** (and is in **free-fall**, e.g., Earth-centered)
- **Modern definition: fixed to quasars** – which exhibit no relative motion on celestial sphere
- **International Celestial Reference Frame (ICRF)** based on coordinates of **295** stable quasars

Newton's Second Law of Motion

- Time-rate of change of **linear momentum** equals **applied force**, F

$$\frac{d}{dt}(m_i \dot{\mathbf{x}}) = F$$



- m_i is the **inertial mass** of the test body $(m_i = \text{constant} \rightarrow m_i \ddot{\mathbf{x}} = F)$

- In the presence of a **gravitational field**, this law must be modified:

$$m_i \ddot{\mathbf{x}} = F + F_g$$

- F_g is a force associated with the **gravitational acceleration** due to a **field** (or space curvature) generated by all masses in the universe, relative to the freely-falling frame (Earth's mass and tidal effects due to moon, sun, etc.)
- action forces, F , and gravitational forces, F_g , are fundamentally different

Newton's Law of Gravitation

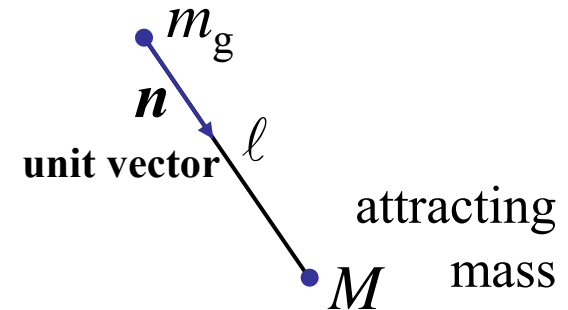


• Gravitational force vector

$$\mathbf{F}_g = G \frac{M m_g}{\ell^2} \mathbf{n} = m_g \mathbf{g}$$

- G = Newton's gravitational constant
- \mathbf{g} = gravitational acceleration due to M
- m_g is the **gravitational mass** of the test body
- it's easier to work with **field potential**, V

$$\mathbf{g} = \nabla V \quad V = GM/\ell$$



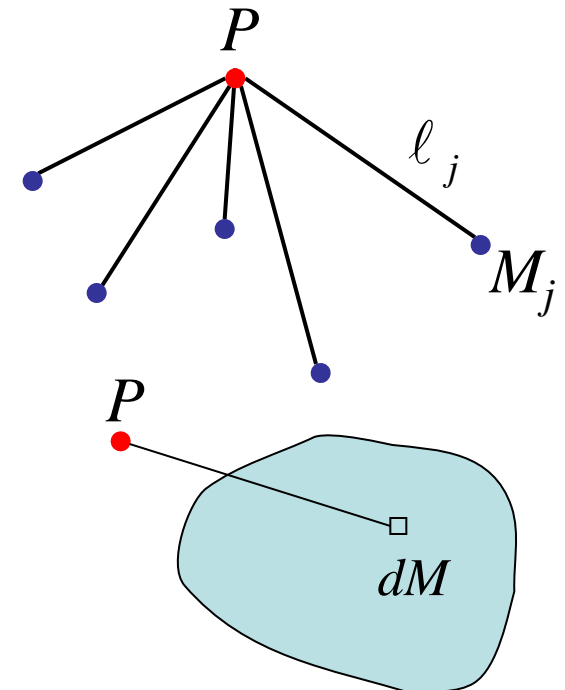
• Many mass points

- law of superposition:

$$V_P = G \sum_j \frac{M_j}{\ell_j}$$

- mass continuum:

$$V_P = G \int_M \frac{dM}{\ell}$$



Equivalence Principle (1)

- **A. Einstein (1907):** No experiment performed in a closed system can distinguish between an **accelerated** reference frame or a reference frame at rest in a uniform **gravitational field**.

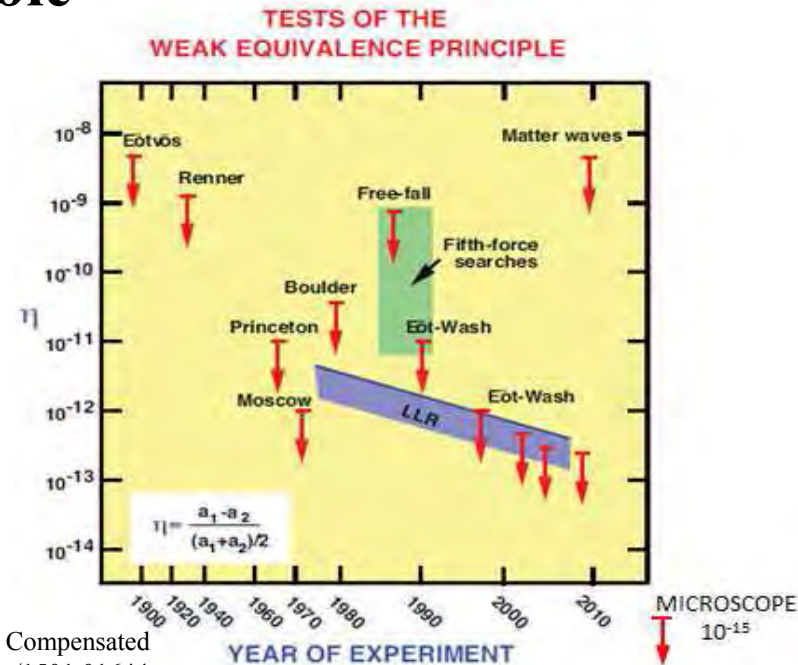
– consequence: inertial mass equals gravitational mass

$$m_i = m_g = m$$

- **Experimental evidence has not been able to dispute this assumption**

– violation of the principle may lead to new theories that unify gravitational and other forces

– proposed French Space Agency mission, **MICROSCOPE***, aims to push the sensitivity by many orders of magnitude



* Micro-Satellite à traînée Compensée pour l'Observation du Principe d'Equivalence (Drag Compensated Micro-satellite to Observe the Equivalence Principle); Berge et al. (2015) <http://arxiv.org/abs/1501.01644>

Equivalence Principle (2)

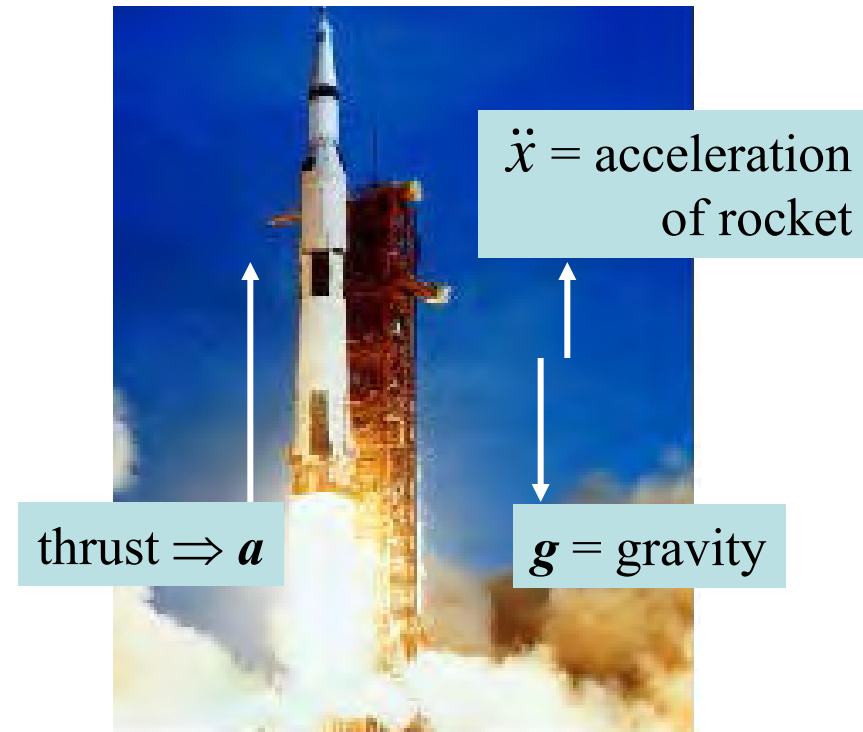
- Equation of motion in the inertial frame

$$\ddot{x}^i = \frac{F^i}{m} + g^i$$

- $\ddot{x} = \frac{d^2 x}{dt^2}$, **vector of total kinematic acceleration**

- $\frac{F^i}{m} = a^i$, **specific force**, or the acceleration resulting from an **action force**; e.g., thrust of a rocket

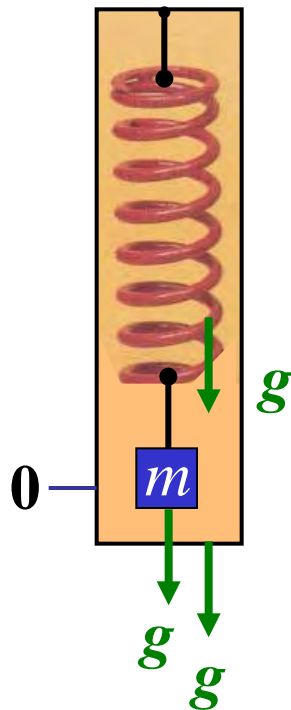
$$\ddot{x}^i = a^i + g^i$$



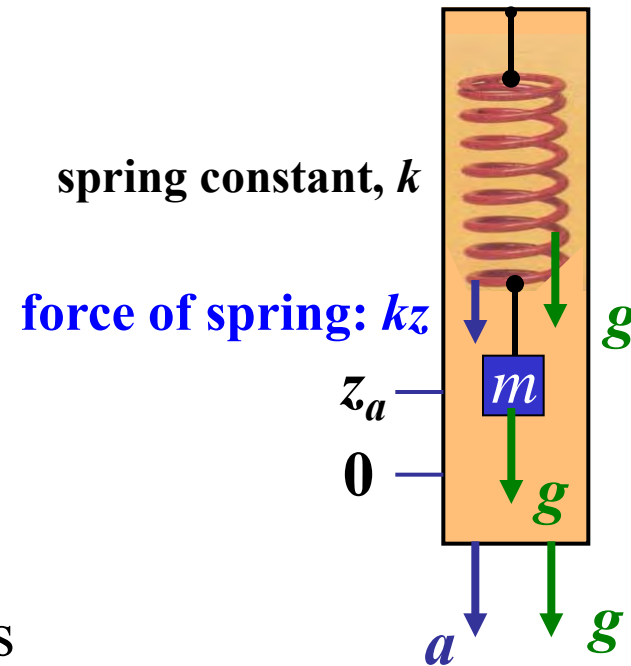
$$|a| \leq |g| \Rightarrow \text{no lift-off !}$$

What Does an Accelerometer Sense?

gravitational field, g
no applied acceleration



gravitational field, g
applied acceleration, a



accelerometer indicates: 0

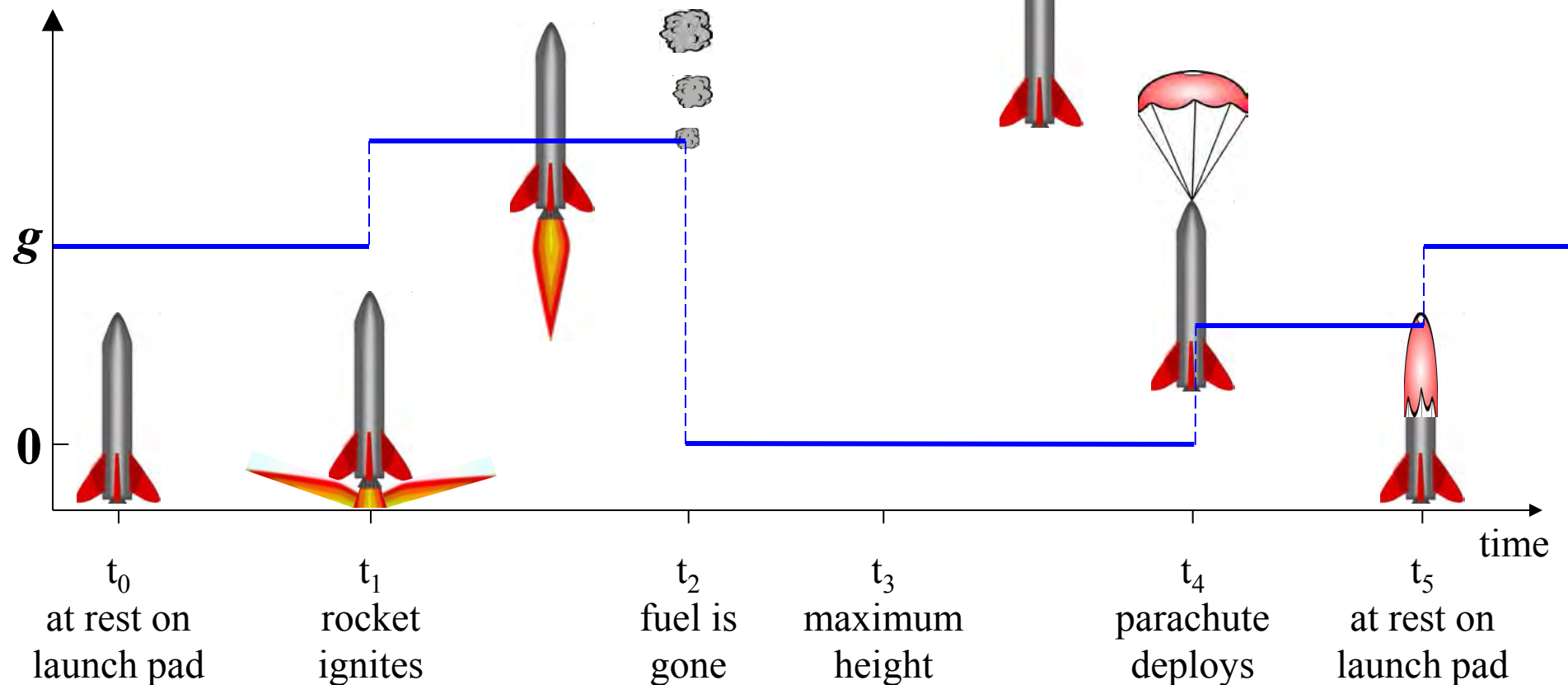
accelerometer indicates: $z_a \sim a$

- Accelerometer does *not* sense gravitation, only acceleration due to **action force** (including *reaction* forces!)

What Does Accelerometer (or Gravimeter) on Rocket Sense?

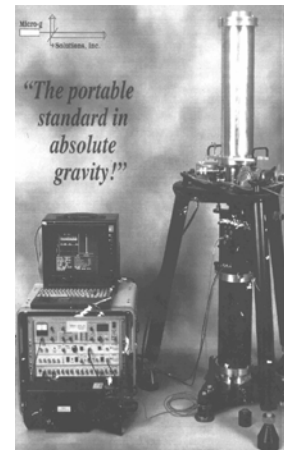
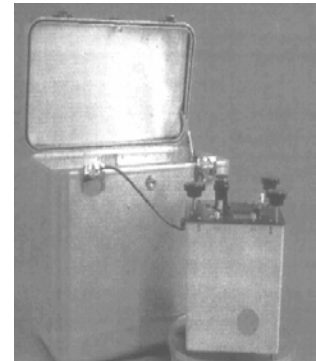
Rocket: experiences no atmospheric drag
engine has constant thrust
launches vertically

Accelerometer axis:
vertically up



Static Gravimetry – Special Case

- Assume non-rotating Earth (for simplicity) $g = \ddot{x} - a$
- Relative (spring) gravimeter: $\ddot{x} = 0 \Rightarrow a = -g$
 - it is an **accelerometer** that senses specific force, a
 - with sensitive axis along **plumb line**, a is the **reaction** force of Earth's surface that keeps the gravimeter from falling
- Absolute (ballistic) gravimeter: $a = 0 \Rightarrow \ddot{x} = g$
 - it **tracks** a test mass in **vacuum** (zero spring force)
 - indirectly, it senses the reaction force that keeps the **reference** from falling
- All operational moving-base gravimeters are **relative sensors**



Basic Equation for Moving-Base Gravimetry

- In the inertial frame:

$$g^i = \ddot{x}^i - a^i$$

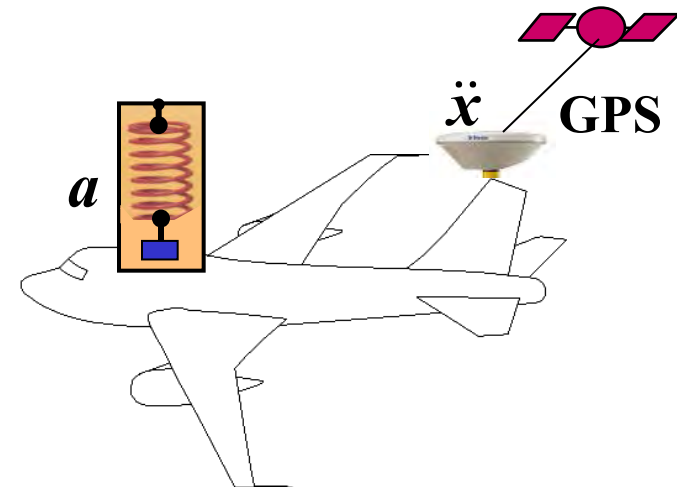
- Because:

- **specific forces** are measured in a non-inertial frame attached to a **rotating body** (vehicle)
- specific forces and kinematic accelerations refer to **different measurement points** of the instrument-carrying vehicle
- generally, gravitation is desired in a **local, Earth-fixed frame**

- Need to introduce:

- **coordinate frames**
- **rotations** and **lever-arm effects**

- Get more complicated expressions for gravimetry equation



Two Possible Approaches to Determine g (1)

- For concepts, consider inertial frame for simplicity: $\ddot{x}^i = a^i + g^i$
- **Position (Tracking) Method** to determine the **unknown**: g

- Integrate equations of motion

$$x^i(t) = x^i(t_0) + \dot{x}^i(t_0)(t - t_0) + \int_{t_0}^t (t - t') (a^i(t') + g^i(t')) dt'$$

- **Positions**, x : from tracking system, like GPS or other GNSS
- **Specific forces**, a : from accelerometer
- method is used for geopotential determination with satellite tracking, and was used also with ground-based inertial positioning systems
- **Advantage**: do not need to differentiate x to get \ddot{x}
- **Disadvantage**: g must be **modeled** in some way to perform the integration (e.g., spherical harmonics in satellite tracking, with statistical constraint)
- Not used for scalar airborne gravimetry due to vertical instability of integral
 - but can be (is) used for horizontal components of gravity!

Two Possible Approaches to Determine g (2)

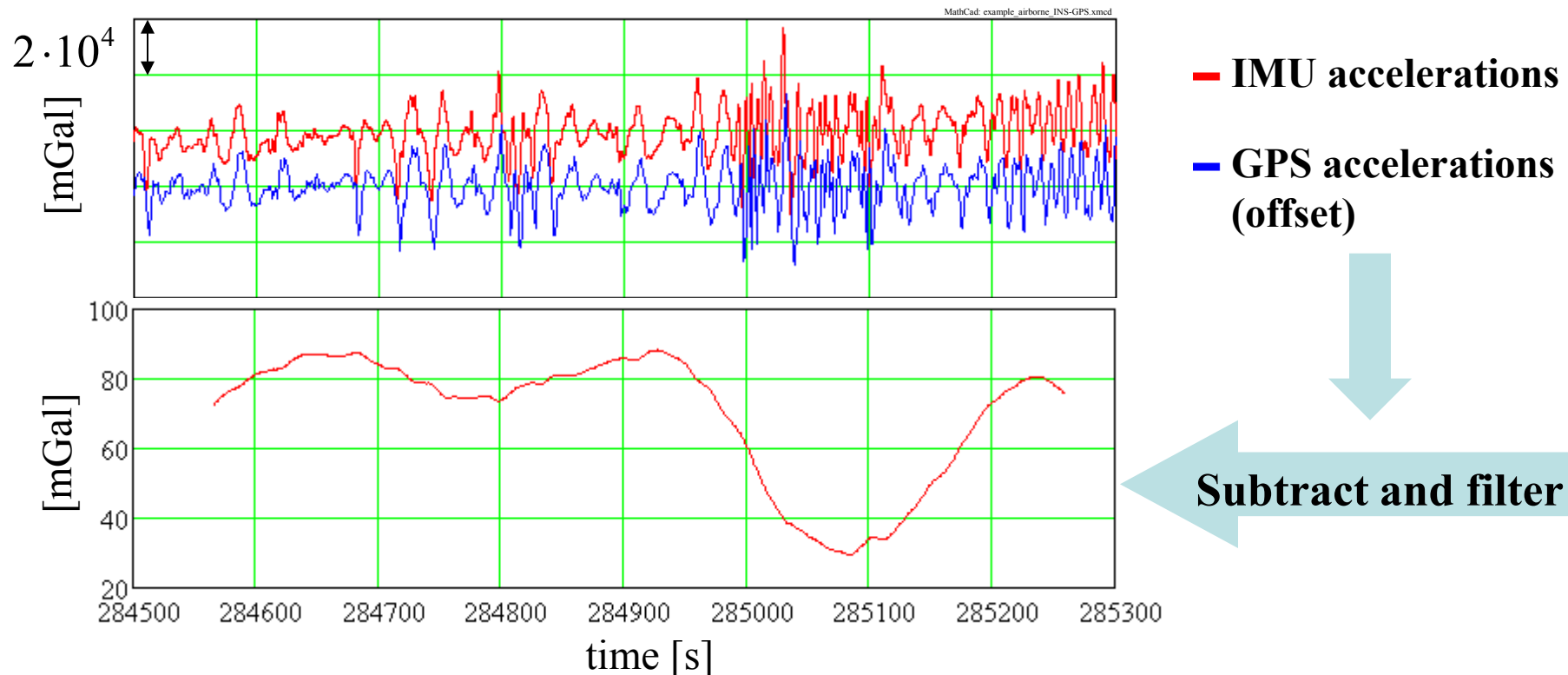
- **Accelerometry Method** to determine the **unknown: g**

$$g^i = \ddot{x}^i - a^i$$

- **Specific force, a** : from accelerometer
 - **Kinematic acceleration, \ddot{x}** : by differentiating position from tracking system, like GPS (GNSS)
 - **Advantage:** g does not need to be modeled
 - **Disadvantage:** positions are processed with two **numerical differentiations**
 - advanced numerical techniques → may be less serious than gravity modeling problem
- **Either position method or accelerometry method requires two independent sensor systems**
 - Tracking system
 - Accelerometer (gravimeter)
 - **Gravimetry accuracy** depends equally on the precision of both systems

The Challenge of Airborne Gravimetry

- Both systems measure large signals e.g., ($> \pm 10000$ mGal)
- Desired gravity disturbance is orders of magnitude smaller
 - **signal-to-noise** ratio may be very small, depending on system accuracies
- e.g., INS/GPS system – data from University of Calgary, 1996



Coordinate Frames

- **Other coordinate frames**
 - **rotating** with respect to inertial frame,
 - may have different **origin point**,
 - have different form of **Newton's law of motion**,
 - all defined by three mutually orthogonal, usually right-handed axes (Cartesian coordinates).

- **Specific frames to be considered:**
 - **navigation frame**: frame in which navigation equations are formulated; usually identified with local North-East-Down (**NED**) directions (***n*-frame**).
 - **Earth-centered-Earth-fixed frame**: frame with origin at Earth's center of mass and axes defined by conventional pole and Greenwich meridian (Cartesian or **geodetic coordinates**) (***e*-frame**).

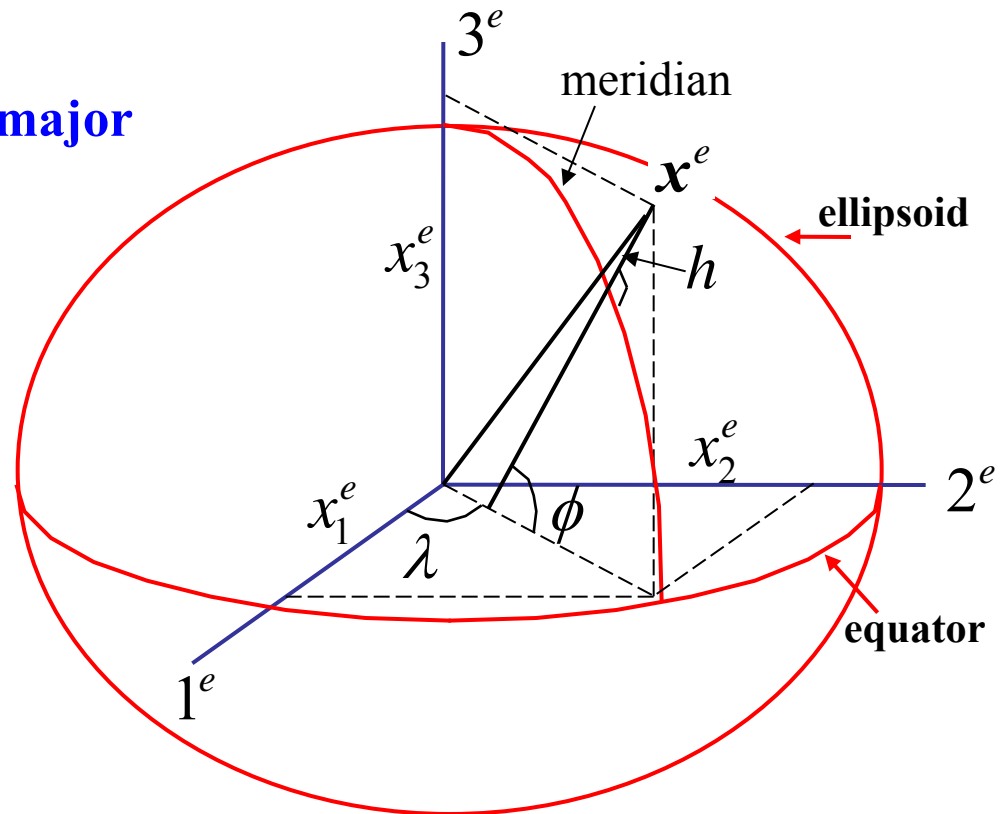
Earth-Centered-Earth-Fixed Coordinates

- Cartesian coordinate vector in e -frame: $\mathbf{x}^e = \begin{pmatrix} x_1^e & x_2^e & x_3^e \end{pmatrix}^T$

- Geodetic coordinates, latitude, longitude, height: ϕ, λ, h

- refer to a particular **ellipsoid** (assume geocentric) with **semi-major axis**, a , and **first eccentricity**, e

- are orthogonal **curvilinear** e -frame coordinates



Transforming Between Cartesian & Geodetic Coordinates

$$x_1^e = (N + h) \cos \phi \cos \lambda$$

$$x_2^e = (N + h) \cos \phi \sin \lambda$$

$$x_3^e = (N(1 - e^2) + h) \sin \phi$$

$$\phi = \tan^{-1} \left(\frac{x_3^e}{\sqrt{(x_1^e)^2 + (x_2^e)^2}} \left(1 + \frac{e^2 N \sin \phi}{x_3^e} \right) \right)$$

$$\lambda = \tan^{-1} (x_2^e / x_1^e)$$

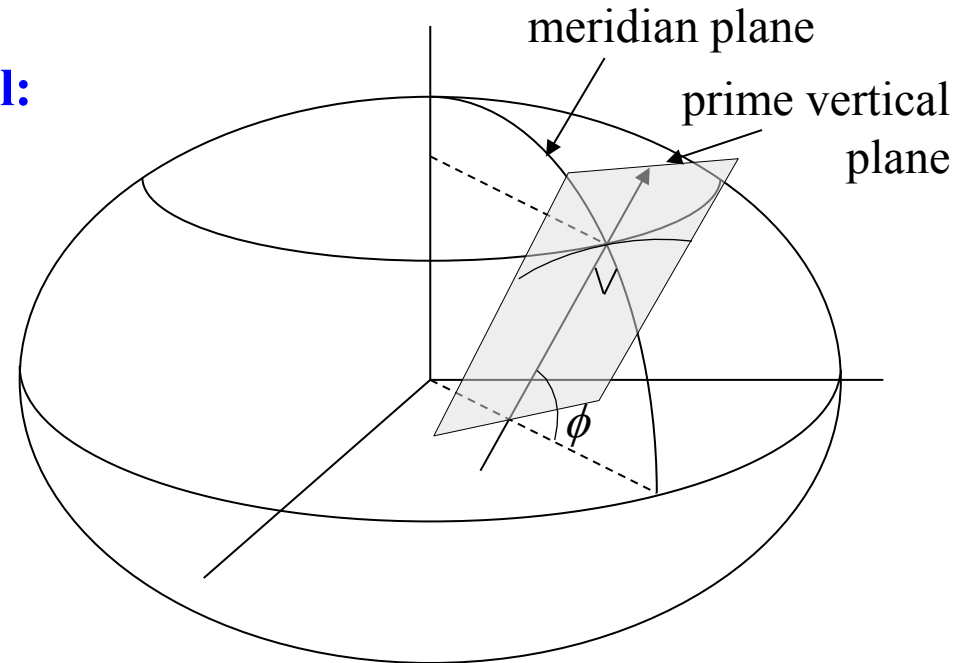
$$h = \sqrt{(x_1^e)^2 + (x_2^e)^2} \cos \phi + x_3^e \sin \phi - a^2 / N$$

– **radius of curvature in prime vertical:**

$$N = \frac{a}{\sqrt{1 - e^2 \sin^2 \phi}}$$

– **radius of curvature in meridian:**

$$M = \frac{a(1 - e^2)}{(1 - e^2 \sin^2 \phi)^{3/2}}$$



• **Cartesian and geodetic coordinates may be used interchangeably**

Rotations and Angular Rates Between Frames

- Assume common origin for frames
- \mathbf{C}_t^s = matrix that rotates coordinates **from** t -frame **to** s -frame
 - vector, \mathbf{x} : $\mathbf{x}^s = \mathbf{C}_t^s \mathbf{x}^t$
 - matrix, \mathbf{A} : $\mathbf{A}^s = \mathbf{C}_t^s \mathbf{A}^t \mathbf{C}_s^t$
 - \mathbf{C}_t^s is **orthogonal**: $\mathbf{C}_s^t \equiv (\mathbf{C}_t^s)^{-1} = (\mathbf{C}_t^s)^T$
- $\boldsymbol{\omega}_{st}^t$ = angular rate vector of t -frame relative to s -frame; components in t -frame

• Let $\boldsymbol{\omega}_{st}^t = \begin{pmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{pmatrix}$ then $[\boldsymbol{\omega}_{st}^t \times] \equiv \boldsymbol{\Omega}_{st}^t = \begin{pmatrix} 0 & -\omega_3 & \omega_2 \\ \omega_3 & 0 & -\omega_1 \\ -\omega_2 & \omega_1 & 0 \end{pmatrix}$

– cross-product is same as multiplication by **skew-symmetric** matrix

• **Time-derivative:** $\dot{\mathbf{C}}_t^s = \mathbf{C}_t^s [\boldsymbol{\omega}_{st}^t \times] = \mathbf{C}_t^s \boldsymbol{\Omega}_{st}^t$

Earth-Fixed vs. Inertial Frames

- Transformation of coordinates between i -frame and e -frame is just a rotation about the 3-axis

$$\mathbf{x}^i = \mathbf{C}_e^i \mathbf{x}^e \Rightarrow \ddot{\mathbf{x}}^i = \ddot{\mathbf{C}}_e^i \mathbf{x}^e + 2\dot{\mathbf{C}}_e^i \dot{\mathbf{x}}^e + \mathbf{C}_e^i \ddot{\mathbf{x}}^e$$

$$\boldsymbol{\omega}_{ie}^e = \begin{pmatrix} 0^* & 0^* & \omega_E \end{pmatrix}^T$$

ω_E = Earth's rotation rate

0^* = neglect rates of polar motion and precession/nutation

$$d\boldsymbol{\omega}_{ie}^e / dt = \mathbf{0}$$

$$\ddot{\mathbf{x}}^i = \mathbf{C}_e^i \left(\boldsymbol{\Omega}_{ie}^e \boldsymbol{\Omega}_{ie}^e + \dot{\boldsymbol{\Omega}}_{ie}^e \right) \mathbf{x}^e + 2\mathbf{C}_e^i \boldsymbol{\Omega}_{ie}^e \dot{\mathbf{x}}^e + \mathbf{C}_e^i \ddot{\mathbf{x}}^e$$

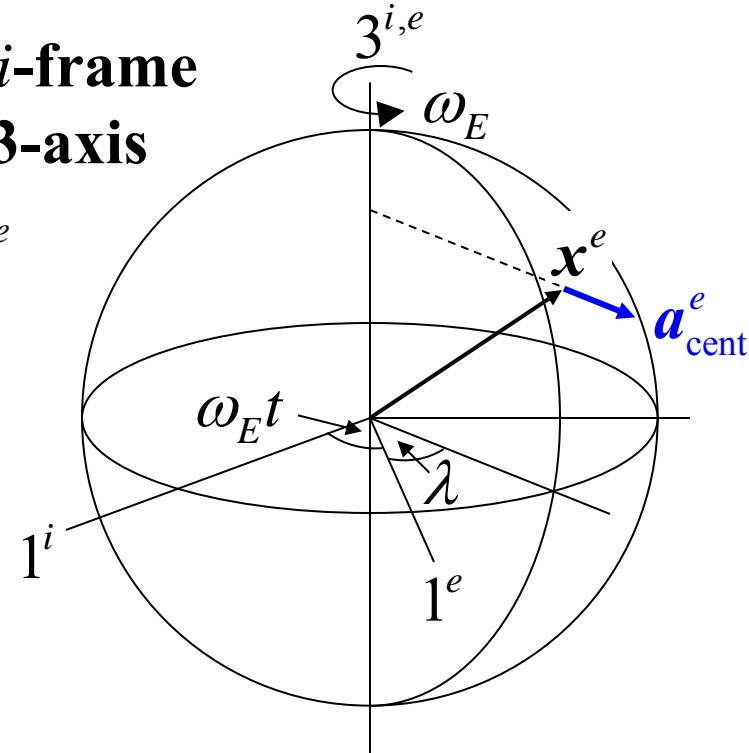
- Extract **centrifugal** acceleration from other kinematic accelerations

$$\mathbf{C}_i^e \ddot{\mathbf{x}}^i = \boldsymbol{\Omega}_{ie}^e \boldsymbol{\Omega}_{ie}^e \mathbf{x}^e + 2\boldsymbol{\Omega}_{ie}^e \dot{\mathbf{x}}^e + \ddot{\mathbf{x}}^e = \mathbf{a}^e + \mathbf{g}^e$$

$$= -\mathbf{a}_{\text{cent}}^e + \mathbf{q}^e$$

← defines \mathbf{q}

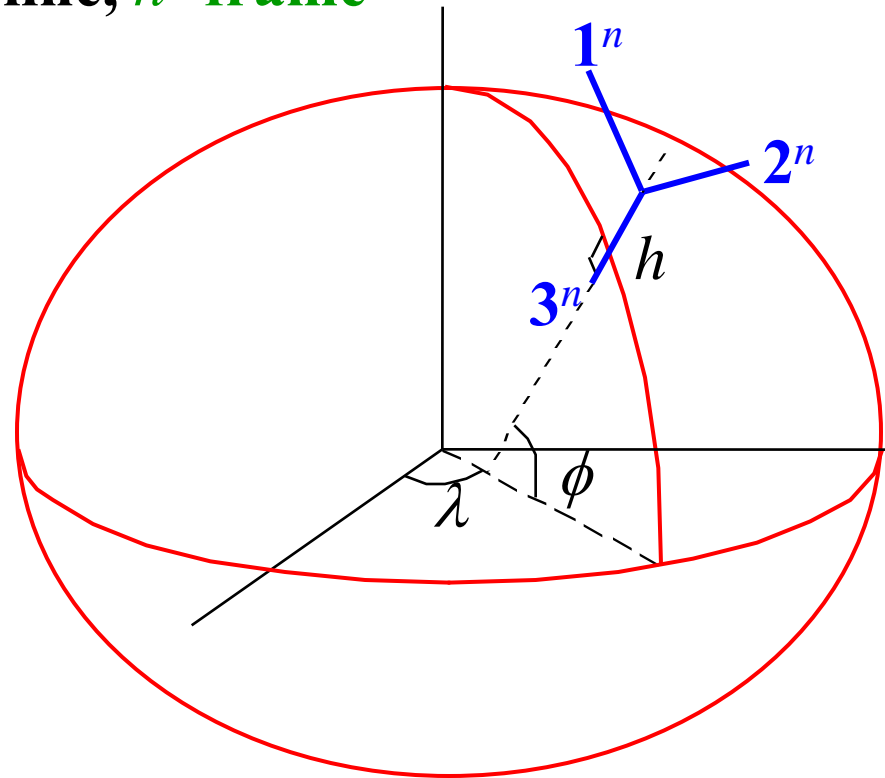
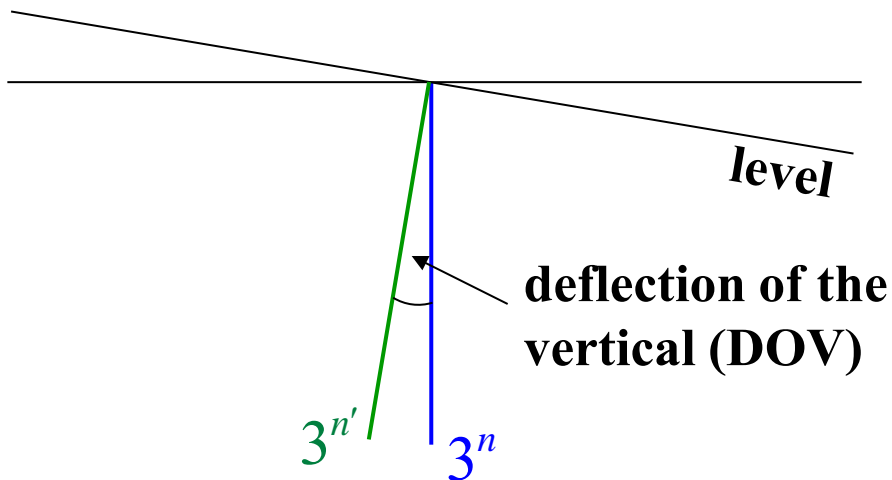
$$\Rightarrow \bar{\mathbf{g}}^e = \mathbf{q}^e - \mathbf{a}^e$$



Navigation Frame

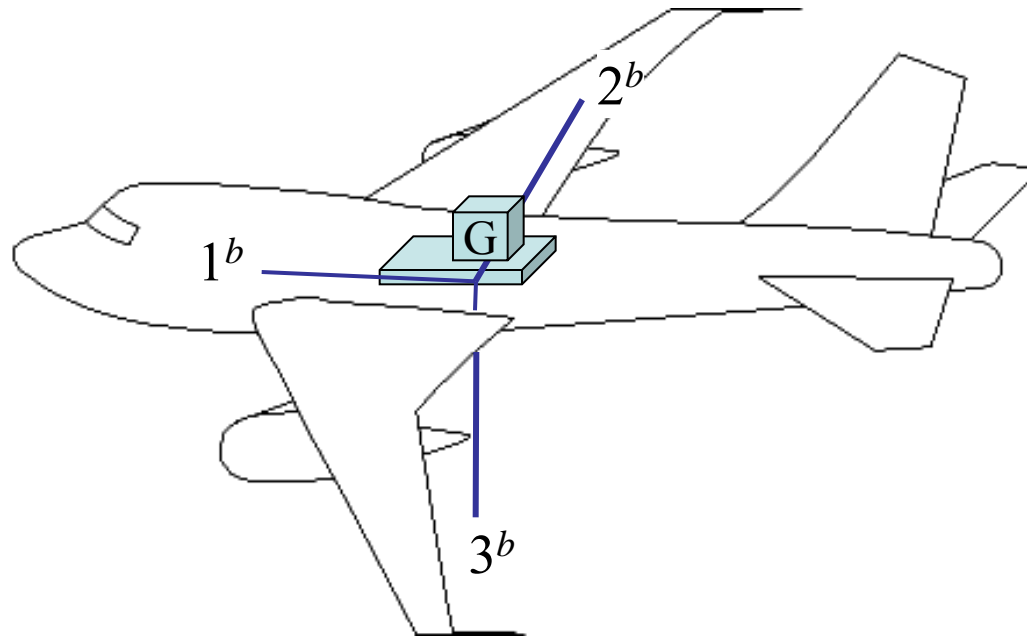
- Usually, north-east-down (**NED**)-frame, or **n -frame**
 - moves with the vehicle – *not* used for coordinates of the vehicle
 - used as reference for **velocity** and **orientation** of the vehicle; and, **gravity**
- **Alternative: vertical along plumb line, n' -frame**

- conventional reference for terrestrial gravity
- no “horizontal” gravity components



Body Frame

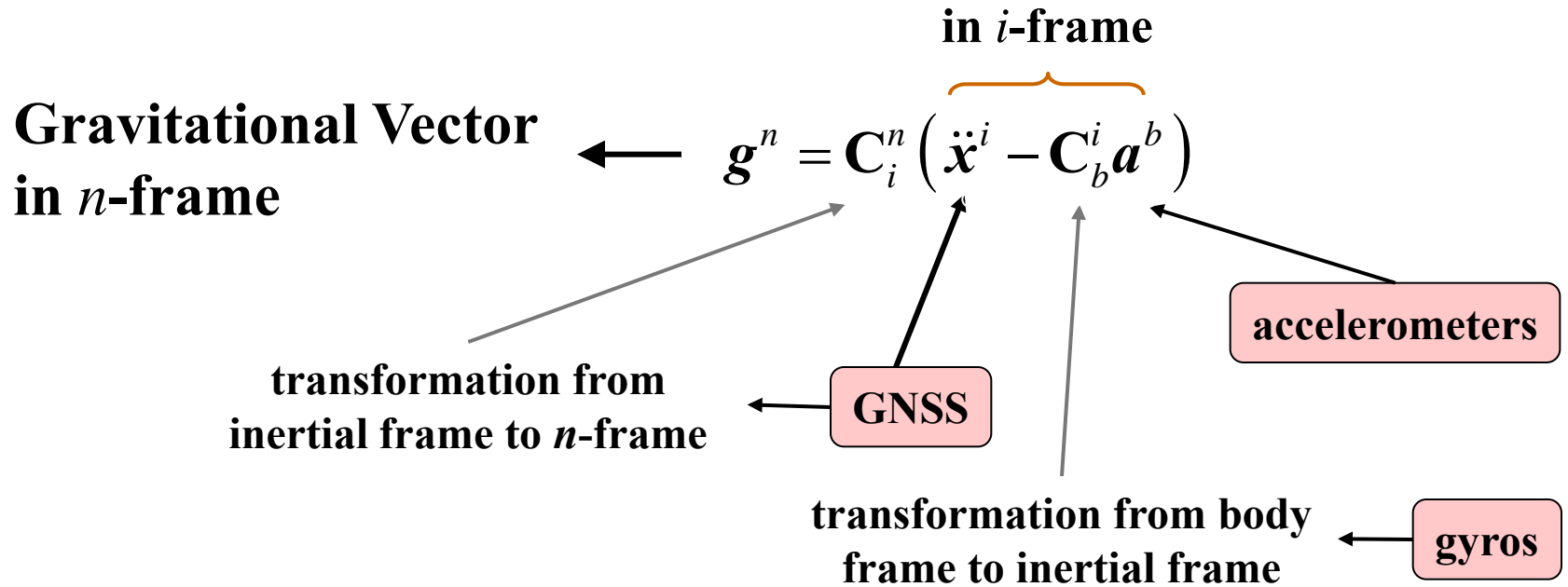
- Axes are defined by principal axes of the vehicle: **forward** (1), **to-the-right** (2), and **through-the-floor** (3)



- **Gravimeter (G) measurements are made either:**
 - in the ***b***-frame – **strapdown** system; gyro data provide orientation
 - in the ***n***-, ***n'***-frames – **platform** is stabilized using IMUs*

* inertial measurement units

Moving-Base Gravimetry – Strapdown Mechanization



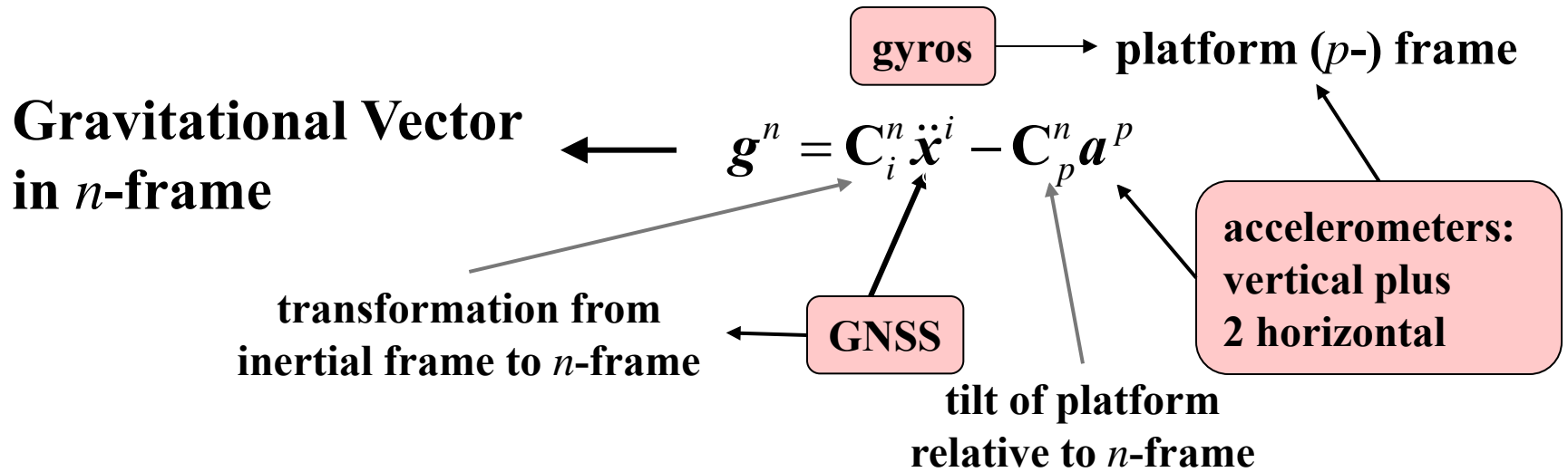
\mathbf{a}^b – **inertial accelerations** measured by accelerometers in **body frame**

$\ddot{\mathbf{x}}^i$ – **kinematic accelerations** obtained from GNSS-derived positions, \mathbf{x} , in **i -frame**

$\mathbf{C}_i^n = \begin{pmatrix} -\sin \phi \cos(\lambda + \omega_E t) & -\sin \phi \sin(\lambda + \omega_E t) & \cos \phi \\ -\sin(\lambda + \omega_E t) & \cos(\lambda + \omega_E t) & 0 \\ -\cos \phi \cos(\lambda + \omega_E t) & -\cos \phi \sin(\lambda + \omega_E t) & -\sin \phi \end{pmatrix}$ – **transformation obtained from GNSS-derived positions, ϕ, λ**

– **lever-arm effects** are assumed to be applied

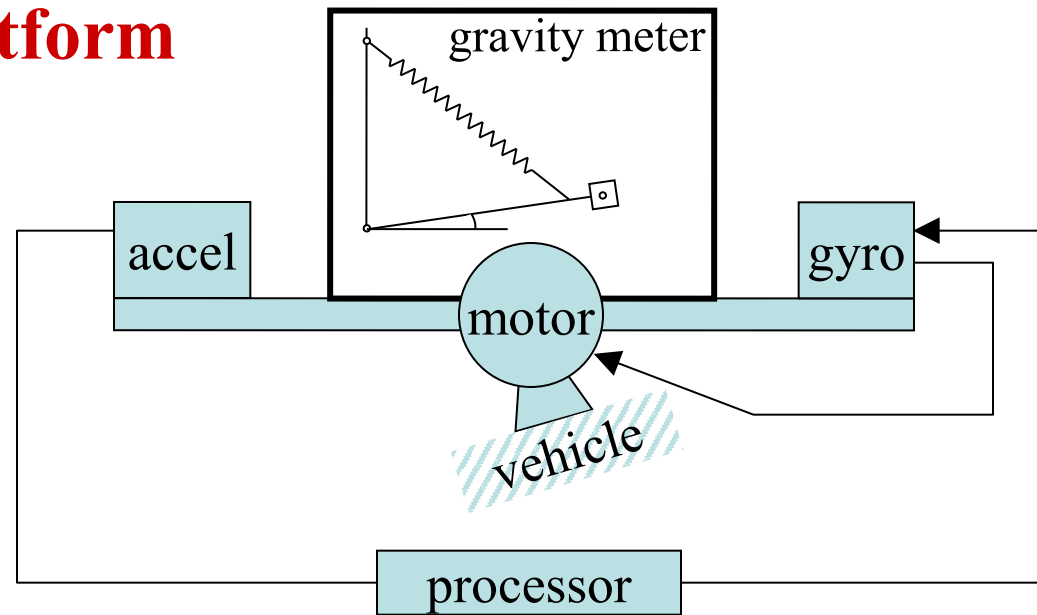
Moving-Base Gravimetry – Stabilized Mechanization



- **Inertial accelerations, \mathbf{a}^p , from accelerometers in platform frame**
- **Mechanizations**
 - **two-axis damped platform** – level (n' -frame) alignment using gyro-driven gimbaled platform in the short term and **mean zero output** of horizontal accelerometers in the long term
 - adequate for benign dynamics
 - **Schuler-tuned inertial stabilized platform** – alignment to n -frame based on inertial /GNSS navigation solution and gyro-driven platform stabilization
 - ideally, $\mathbf{C}_p^n = \mathbf{I}$; but note, n -frame differs from n' -frame by deflection of the vertical
 - better for more dynamic environments

Two-Axis Stabilized Platform

- Schematic for one axis



- **horizontal** accelerometer, through processor, ensures that gyroscope reference direction is **precessed** to account for Earth rotation and curvature

- zero acceleration implies level orientation (**without** horizontal specific forces!)

- ad hoc damping of platform by processor

- corrects gyro drift, but is subject to accelerometer bias

- **Schuler-tuned** three-axis stabilization: more accurate IMUs and n -frame stabilization (using navigation solution velocity in n -frame)

Scalar Moving-Base Gravimetry

- Determine the **magnitude of gravity** – the **plumb line** component
 - consistent with ground-based measurement (recall gravimeter is **leveled**)

- gravitation vector: $\mathbf{g}^n = \mathbf{C}_i^n \ddot{\mathbf{x}}^i - \mathbf{C}_b^n \mathbf{a}^b = \left(-\mathbf{a}_{\text{cent}}^n + \mathbf{q}^n \right) - \mathbf{a}^n$

- **gravity** vector: $\bar{\mathbf{g}}^n = \mathbf{g}^n + \mathbf{a}_{\text{cent}}^n$

$$\boxed{\bar{\mathbf{g}}^n = \mathbf{q}^n - \mathbf{a}^n} \quad \text{GNSS} \rightarrow \mathbf{q}^n$$

- this holds in **any frame!** e.g., $\bar{\mathbf{g}}^{n'} = \mathbf{q}^{n'} - \mathbf{a}^{n'}$

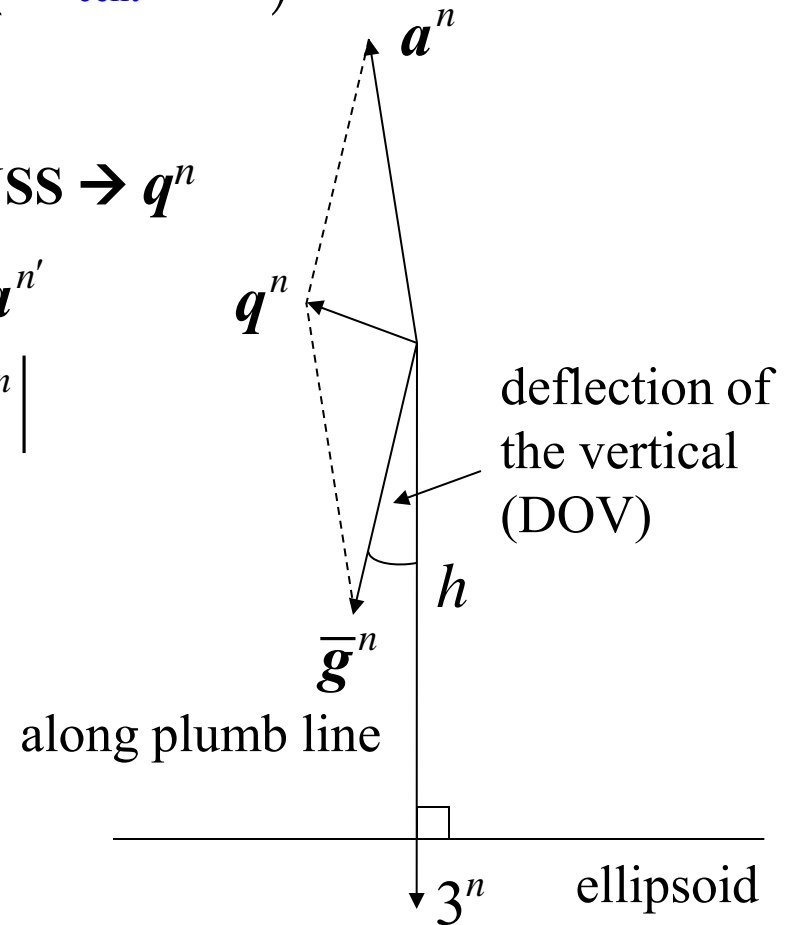
- note: straight and level flight $\rightarrow |\mathbf{q}^n| \square |\mathbf{a}^n|$

- thus: $\bar{g} = |\bar{\mathbf{g}}^n|$, but note: $\bar{g} \neq \bar{g}_3$

- n -frame mechanization does **not** account for the **deflection of the vertical**

- since n' -frame is aligned to plumb line,

$$\bar{g} = \bar{g}_3^{n'}$$



Unconstrained Scalar Gravimetry

- **One Option:** $\bar{g} = \left| \bar{\mathbf{g}}^n \right| = \left| \mathbf{q}^n - \mathbf{a}^n \right|$
 - **unconstrained** in the sense that the frame for vectors is arbitrary (n -frame is used for illustration)
 - also known as **strapdown inertial scalar gravimetry (SISG)**
 - $\mathbf{q}^n = \mathbf{C}_i^n \ddot{\mathbf{x}}^i + \mathbf{a}_{\text{cent}}^n$ obtained exclusively from GNSS
 - $\mathbf{a}^n = \mathbf{C}_b^n \mathbf{a}^b$ requires **orientation** of b -frame (relative to n -frame)
- **Requires comparable accuracy in all accelerometers and precision gyros if platform is arbitrary (e.g., strapdown)**
- **Calgary group demonstrated good results (e.g., Glennie and Schwarz 1999); see also (Czompo and Ferguson 1995)**

Rotation-Invariant Scalar Gravimetry (RISG)

- Another option to get \bar{g} : based on **total** specific force from gravimeter and orthogonal accelerometers

$$a^2 = (a_1^p)^2 + (a_2^p)^2 + (a_3^p)^2 \quad \left(= (a_1^{n'})^2 + (a_2^{n'})^2 + (a_3^{n'})^2 \right)$$

- Then

$$a_3^{n'} = \sqrt{a^2 - (a_1^{n'})^2 - (a_2^{n'})^2}$$

$$= \sqrt{a^2 - (q_1^{n'} - \bar{g}_1^{n'})^2 - (q_2^{n'} - \bar{g}_2^{n'})^2}$$

$$= \sqrt{a^2 - (q_1^{n'})^2 - (q_2^{n'})^2} \quad \text{since } \bar{g}_1^{n'} = 0 = \bar{g}_2^{n'}$$

$$\bar{g} \equiv \bar{g}_3^{n'} \approx q_3^n - \sqrt{a^2 - (q_1^n)^2 - (q_2^n)^2} \quad \text{neglecting the DOV, } q^{n'} = q^n$$

- **Platform orientation** is not specifically needed for a^2
 - however, errors in q^n , being squared, tend to **bias** the result (Olesen 2003)

RISG Approach in More Detail (1)

- Define **Earth-fixed** velocity vector in the n -frame

$$\mathbf{v}^n = \mathbf{C}_e^n \dot{\mathbf{x}}^e = \begin{pmatrix} v_N \\ v_E \\ v_D \end{pmatrix} = \begin{pmatrix} \dot{\phi}(M+h) \\ \dot{\lambda}(N+h)\cos\phi \\ -\dot{h} \end{pmatrix} \quad \mathbf{C}_e^n = \begin{pmatrix} -\sin\phi\cos\lambda & -\sin\phi\sin\lambda & \cos\phi \\ -\sin\lambda & \cos\lambda & 0 \\ -\cos\phi\cos\lambda & -\cos\phi\sin\lambda & -\sin\phi \end{pmatrix}$$

- It can be shown (Appendix A) that

$$\mathbf{q}^n = \frac{d\mathbf{v}^n}{dt} + (\boldsymbol{\Omega}_{ie}^n + \boldsymbol{\Omega}_{in}^n) \mathbf{v}^n \quad \boldsymbol{\Omega}_{ie}^n + \boldsymbol{\Omega}_{in}^n = \begin{pmatrix} 0 & (\dot{\lambda} + 2\omega_e)\sin\phi & -\dot{\phi} \\ -(\dot{\lambda} + 2\omega_e)\sin\phi & 0 & -(\dot{\lambda} + 2\omega_e)\cos\phi \\ \dot{\phi} & (\dot{\lambda} + 2\omega_e)\cos\phi & 0 \end{pmatrix}$$

- Thus, strictly from GNSS, the third component is

$$q_3^n = -\ddot{h} + 2\omega_e v_E \cos\phi + \frac{v_N^2}{M+h} + \frac{v_E^2}{N+h}$$

RISG Approach in More Detail (2)

- Inertial acceleration includes **tilt error** if platform is not level

$$a_3^{n'} = \left(\mathbf{C}_p^{n'} \mathbf{a}^p \right)_3 = a_3^p - \delta a_{\text{tilt}} \quad \left(= \sqrt{a^2 - (q_1^{n'})^2 - (q_2^{n'})^2} \right)$$

- Kinematic acceleration (from GNSS) includes neglect of **DOV**

$$q_3^{n'} = \left(\mathbf{C}_n^{n'} \mathbf{q}^n \right)_3 = q_3^n - \delta q_{\text{DOV}}$$

- Third component of $\bar{\mathbf{g}}^{n'} = \mathbf{q}^{n'} - \mathbf{a}^{n'}$, **along plumb line**,

$$\bar{g} = \underset{\substack{\uparrow \\ \text{gravimeter}}}{-a_3^p} - \ddot{h} + \underbrace{2\omega_E v_E \cos \phi + \frac{v_N^2}{M+h} + \frac{v_E^2}{(N+h)}}_{\delta g_{\text{Eötvös}}} + \delta a_{\text{tilt}} - \delta q_{\text{DOV}}$$

Eötvös Effect



Loránd Eötvös
1848 - 1919

- **Exact in n -frame (note: $v_{N,E}$ at altitude!)**

$$\delta g_{\text{Eötvös}} = 2\omega_E v_E \cos \phi + \frac{v_N^2}{M+h} + \frac{v_E^2}{(N+h)}$$

- **Approximations**

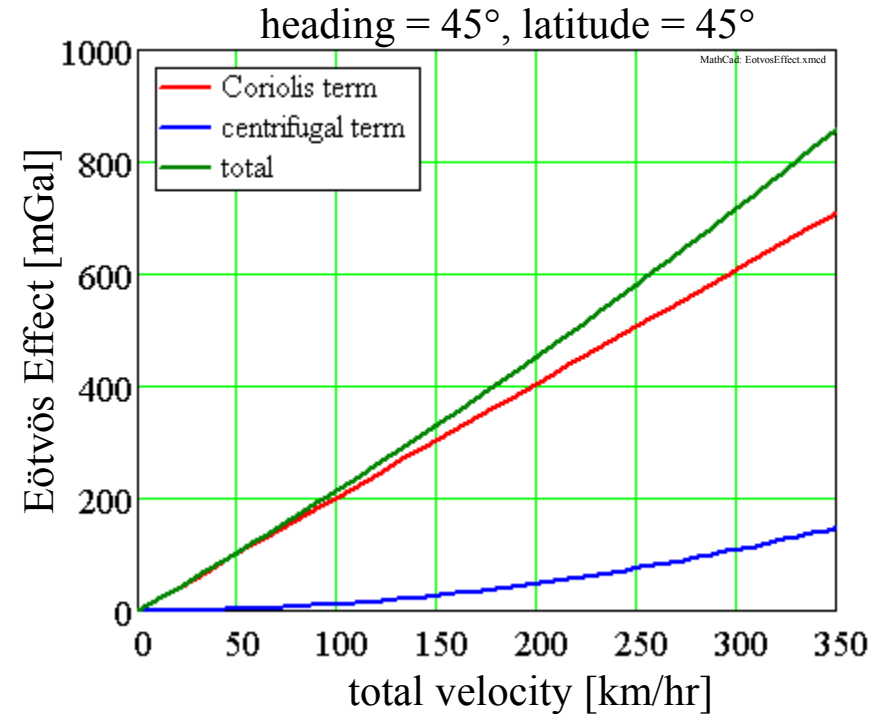
– **spherical:**

$$\delta g_{\text{Eötvös}} \approx 2\omega_E v_E \cos \phi + \frac{v^2}{R+h}$$

– **first-order ellipsoidal (Harlan 1968):**

$$\delta g_{\text{Eötvös}} \approx \frac{v^2}{a} \left(1 + \frac{h}{a} - f \left(1 - \cos^2 \phi (3 - 2 \sin^2 \alpha) \right) \right) + 2v\omega_E \sin \alpha \cos \phi \left(1 + \frac{h}{a} \right) + O(f^2)$$

a = ellipsoid semi-major axis; α = azimuth; v = **ground speed!**

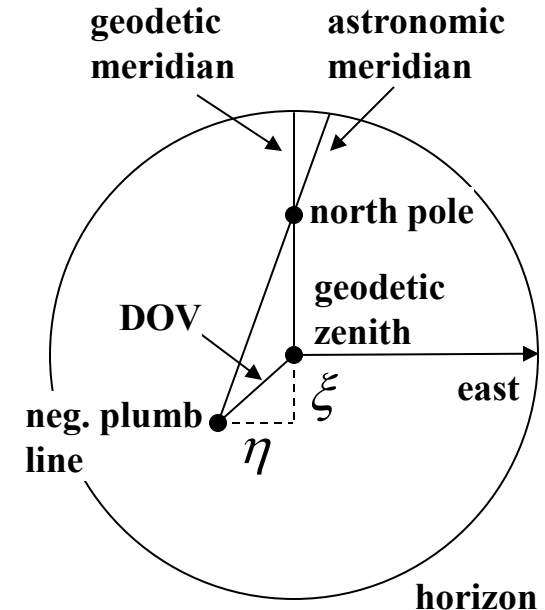


DOV Error in Kinematic Acceleration

- DOV components define the **small** angles between the n - and n' -frames

$$\mathbf{C}_n^{n'} = \mathbf{R}_1(-\eta)\mathbf{R}_2(\xi) = \begin{pmatrix} 1 & 0 & -\xi \\ 0 & 1 & -\eta \\ \xi & \eta & 1 \end{pmatrix}$$

– ignore rotation about 3-axis



- **DOV error**

$$\delta q_{\text{DOV}} = q_3^n - (\mathbf{C}_n^{n'} \mathbf{q}^n)_3 = -\xi q_1^n - \eta q_2^n$$

- Assume $\text{rms}(\text{DOV}) = 10 \text{ arcsec}$, $q_{1,2} = 10^4 \text{ mGal}$

$$\text{rms}(\delta q_{\text{DOV}}) = 0.7 \text{ mGal}$$

- This error is correctable, e.g., using **EGM2008** deflection model

Tilt Error (1)

- One way to compute tilt error (p. 3.27)

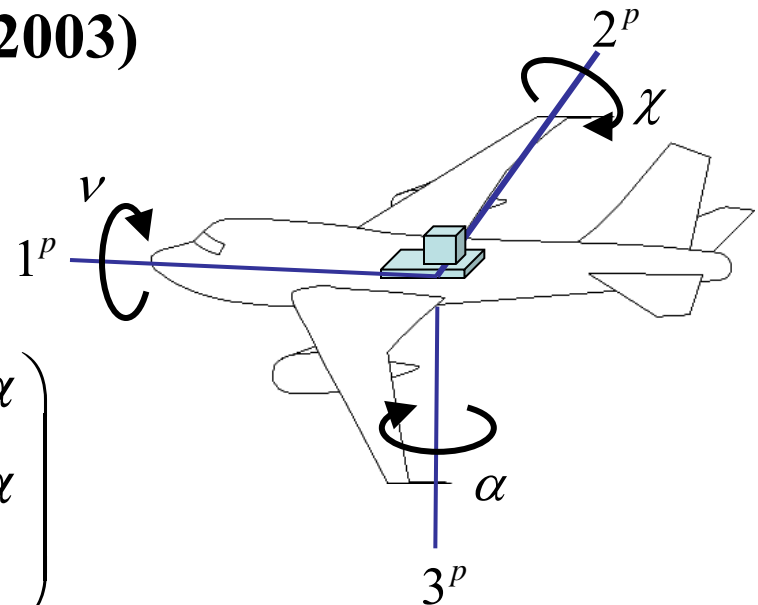
$$\delta a_{\text{tilt}} = a_3^p - \sqrt{a^2 - (q_1^n)^2 - (q_2^n)^2} \quad (\text{neglecting DOV})$$

- random errors in q_1^n, q_2^n are squared and can cause bias (**rectification error**)

- Better model for the tilt error (Olesen 2003)

- define **orientation angles**, ν, χ, α
- assume ν, χ are **small**; α is arbitrary

$$\mathbf{C}_p^n \approx \begin{pmatrix} \cos \alpha & -\sin \alpha & \chi \cos \alpha + \nu \sin \alpha \\ \sin \alpha & \cos \alpha & \chi \sin \alpha - \nu \cos \alpha \\ -\chi & \nu & 1 \end{pmatrix}$$



- thus, approximate platform **stabilization** is required!

Tilt Error (2)

- From $\bar{\mathbf{g}}^n = \mathbf{q}^n - \mathbf{C}_p^n \mathbf{a}^p$,

$$\begin{pmatrix} \bar{g}_1^n \\ \bar{g}_2^n \end{pmatrix} = \begin{pmatrix} q_1^n - a_1^p \cos \alpha + a_2^p \sin \alpha - (\chi \cos \alpha + \nu \sin \alpha) a_3^p \\ q_2^n - a_1^p \sin \alpha - a_2^p \cos \alpha - (\chi \sin \alpha - \nu \cos \alpha) a_3^p \end{pmatrix}$$

- If $\bar{g}_1^n = 0 = \bar{g}_2^n$ (neglecting DOV is second-order effect on tilt error)

$$\chi = \frac{1}{a_3^p} (q_1^n \cos \alpha + q_2^n \sin \alpha - a_1^p)$$

$$\nu = \frac{1}{a_3^p} (q_1^n \sin \alpha - q_2^n \cos \alpha + a_2^p)$$

Tilt angles are **computed**
from accelerometers, GNSS,
and azimuth

- **Third** component of $\delta \mathbf{a}_{\text{tilt}} = \mathbf{a}^p - \mathbf{C}_p^n \mathbf{a}^p$

$$\delta a_{\text{tilt}} \approx a_1^p \chi - a_2^p \nu \quad \text{(first-order approximation)}$$

Tilt Error (3)

- **Tilt error can be written as**

$$\delta a_{\text{tilt}} \approx \frac{q_1^w - a_1^p}{a_3^p} a_1^p + \frac{q_2^w - a_2^p}{a_3^p} a_2^p$$

$$q_1^w = q_1^n \cos \alpha + q_2^n \sin \alpha$$

$$q_2^w = -q_1^n \sin \alpha + q_2^n \cos \alpha$$

– where q_1^w, q_2^w are kinematic accelerations in the “**wander-azimuth**” frame

- **Taking differentials of the specific force components,**

$$\delta(\delta a_{\text{tilt}}) = \frac{(q_1^w - a_1^p) - a_1^p}{a_3^p} \delta a_1^p + \frac{(q_2^w - a_2^p) - a_2^p}{a_3^p} \delta a_2^p - \frac{(q_1^w - a_1^p) a_1^p + (q_2^w - a_2^p) a_2^p}{(a_3^p)^2} \delta a_3^p$$

- **Assume** $(q_{1,2}^w - a_{1,2}^p) = a_3^p \cdot O(\nu, \chi) \square 10^4 \text{ mGal}$, $a_3^p \approx 10^6 \text{ mGal}$, $a_{1,2}^p = O(10^4 \text{ mGal})$

– error in vertical accelerometer (gravimeter) is second-order for tilt correction

– error in horizontal accels. can be 100 worse than tilt correction accuracy

- **In practice, tilt correction is subjected to appropriate filters; see (Olesen 2003)**

Lever-Arm Effect

- Apply to **kinematic acceleration** derived from GNSS tracking
 - assume gravimeter is at center of mass of vehicle
- **In the inertial frame:** $\mathbf{x}_{\text{antenna}}^i = \mathbf{x}_{\text{gravimeter}}^i + \mathbf{b}^i$
 - where $\mathbf{b}^i = \mathbf{C}_b^i \mathbf{b}^b$, $\mathbf{b}^b =$ fixed antenna offset relative to gravimeter
 - \mathbf{C}_b^i – obtained from gyro data
- **Numerical differentiation:** $\ddot{\mathbf{x}}_{\text{gravimeter}}^i = \frac{d^2}{dt^2} (\mathbf{x}_{\text{antenna}}^i - \mathbf{b}^i)$
 - extract relevant component in particular frame
 - n -frame: vertical component of $\mathbf{C}_i^n \ddot{\mathbf{x}}_{\text{gravimeter}}^i \rightarrow \ddot{h}$

Scalar Gravimetry Equation

- Final equation for the **gravity anomaly** at altitude point, P'_a

$$\Delta g_{P'_a} = f - \ddot{h} + \delta g_{\text{Eötvös}} + \delta a_{\text{tilt}} - \delta q_{\text{DOV}} - \gamma_{Q'_a} - (f_0 - g_0)$$

- where $f = -a_3^p$ is the gravity meter reading
- where $\gamma_{Q'_a}$ is **normal gravity** at the **normal height** of P'_a above the ellipsoid
- where $f_0 - g_0$ is the initial offset of the gravimeter reading from true gravity
- gravimeter measurement, f , includes various inherent instrument corrections
- tilt error depends on accuracy of platform accelerometers
- exact and sufficiently approximate formulas exist for normal gravity at Q'_a
- accuracy in \ddot{h} must be **commensurate** with gravimeter accuracy



IV. Overview of Airborne Gravimetry Systems

- **LaCoste/Romberg sea-air gravimeters**
- **Other Airborne gravimeters**

Airborne Gravimeters

- **Instrumentation overview of scalar airborne gravimetry**
 - **LaCoste-Romberg** instruments dominate the field
 - many other instrument types in operation or being tested
- **All gravimeters are single-axis accelerometers**
 - mechanical spring accelerometers (vertical spring, horizontal beam)
 - manual or automatic (**force-rebalance**) nulling
 - torsion wire (horizontal beam, no nulling)
 - electromagnetic spring (force-rebalance)
 - vibrating string accelerometers

Lucien J.B. LaCoste (1908-1995)

“The gravity meters Lucien B. LaCoste invented revolutionized geodesy and gave scientists the ability to precisely measure variations in Earth's gravity from land, water, and space” J.C. Harrison (1996)*



At University of Texas, Austin

- **Inventor**
- **Scientist**
- **Teacher**
- **Entrepreneur**

*https://web.archive.org/web/20080527061634/http://www.agu.org/sci_soc/lacoste.html
Earth in Space Vol. 8, No. 9, May 1996, pp. 12-13. © 1996 American Geophysical Union;
see also (Harrison 1995).

LaCoste-Romberg Air-Sea Model S Gravimeter

- Beam is in **equilibrium** if $\text{torque}(\text{spring}) = \text{torque}(mg)$

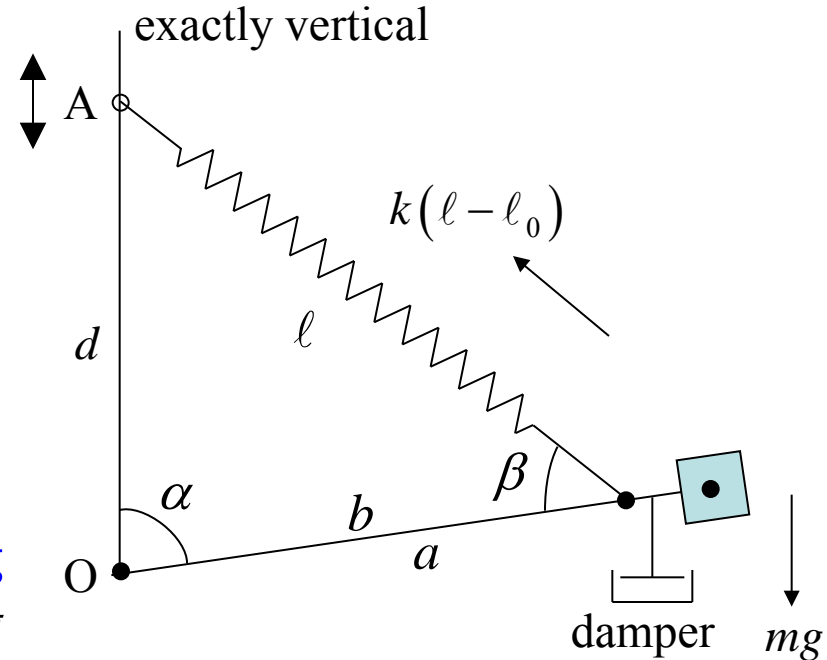
$$k(\ell - \ell_0)b \sin \beta = mga \sin \alpha$$

- law of sines: $\ell \sin \beta = d \sin \alpha$

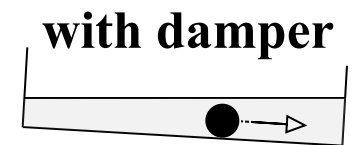
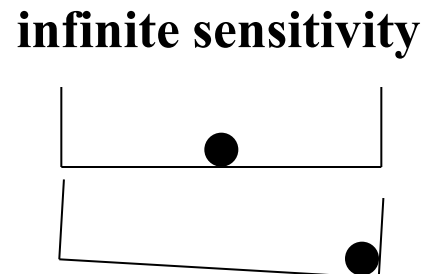
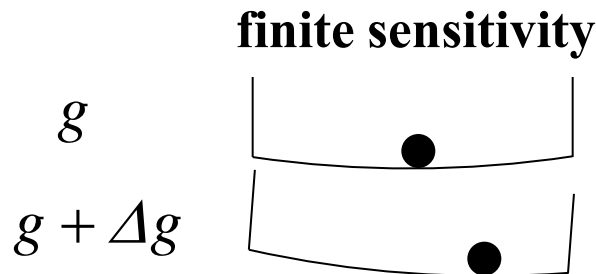
- zero-length spring: $\ell_0 = 0$

$$\Rightarrow kbd = mga$$

- independent of $\alpha \rightarrow$ equilibrium at **any beam position** for a given g
- independent of $\ell \rightarrow$ **no change in spring length** could accommodate a change in g



- From (Valliant 1992):



\rightarrow measure beam **velocity!**

(Inherent) Cross-Coupling Effect

- **Horizontal** accelerations couple into the vertical movement of horizontal beam gravimeters that are **not nulled**

- total torque on beam due to external accelerations:

$$T = ma(\ddot{x} \sin \theta + (g + \ddot{z}) \cos \theta)$$

- it can be shown (LaCoste and Harrison 1961) that the **cross-coupling error** is

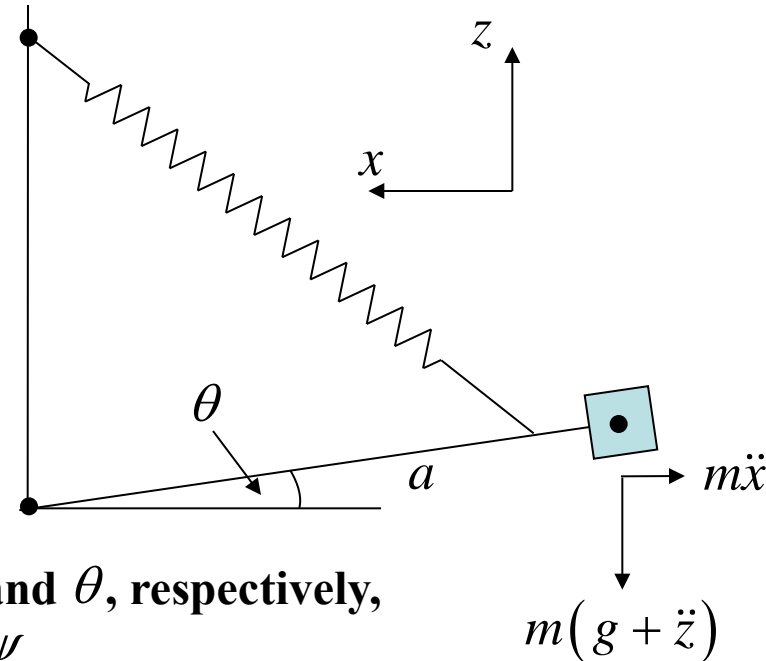
$$\varepsilon = \frac{1}{2} \ddot{x}_1 \theta_1 \cos \psi$$

where \ddot{x}_1, θ_1 are amplitudes of components of \ddot{x} and θ , respectively, that have the **same period and phase difference**, ψ

- e.g., $\theta_1 = 1^\circ$, $\ddot{x}_1 = 0.1 \text{ m/s}^2 \Rightarrow \varepsilon = 90 \text{ mGal}$

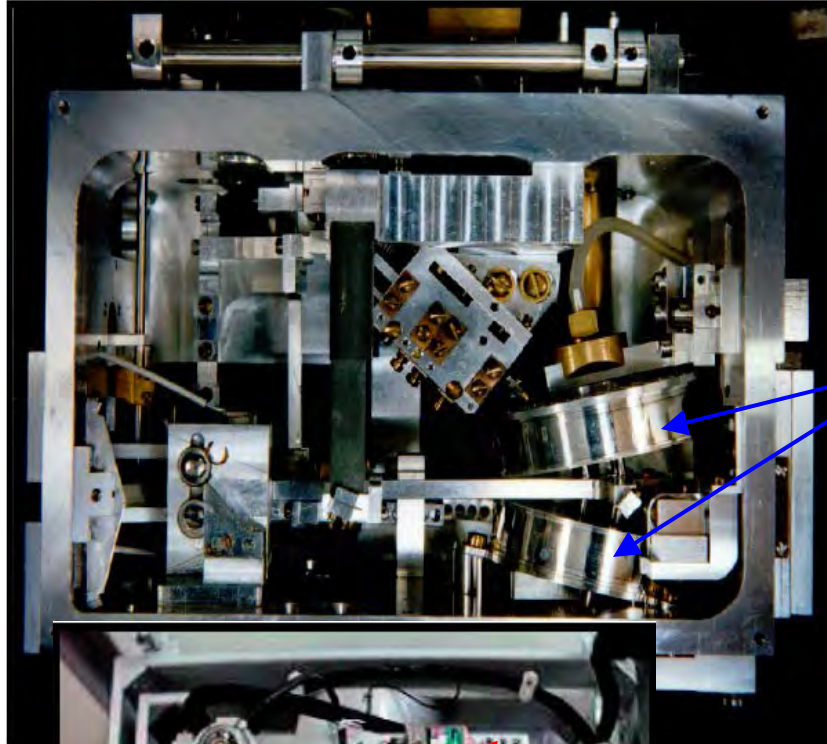
- There is **no** cross-coupling effect for

- force-rebalance gravimeters
- vertical-spring gravimeters



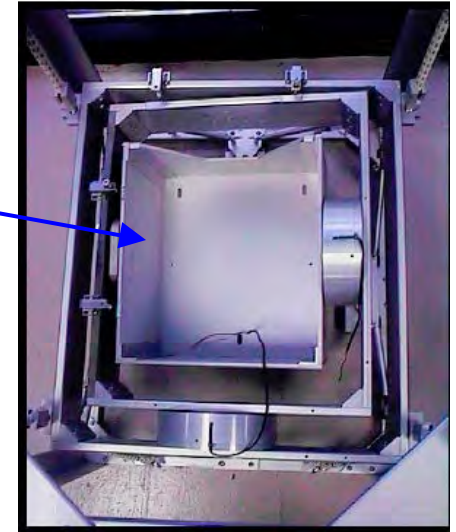
LaCoste-Romberg Model S Sensor and Platform

Interior Side View

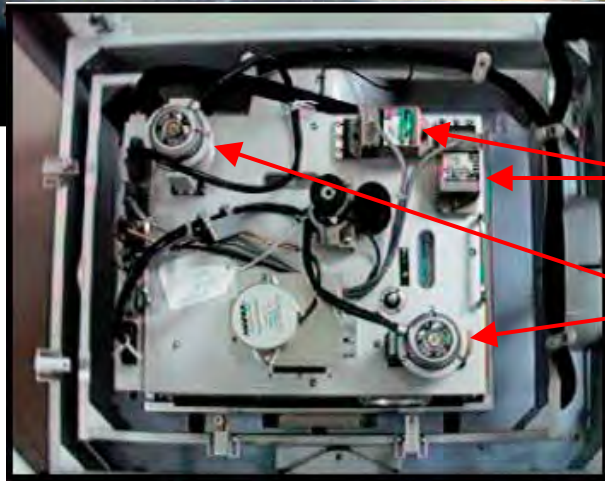


Dampers

Stabilized Platform



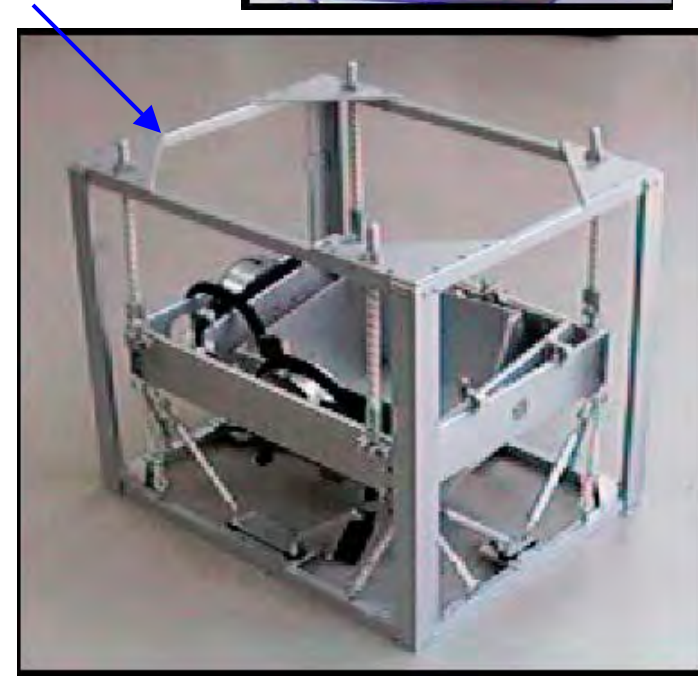
Outer Frame



Accelerometers

Gyroscopes

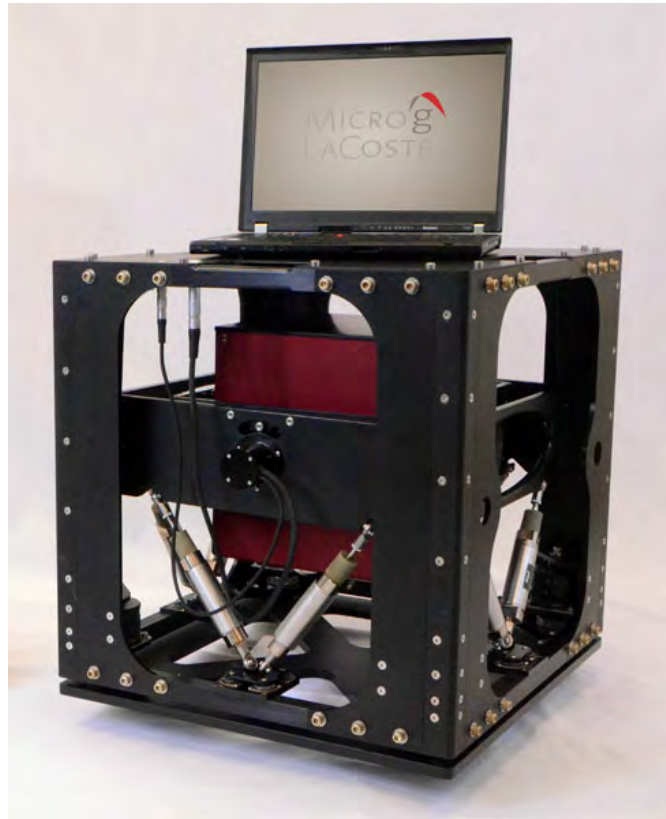
view of top lid



From: Instruction Manual, LaCoste and Romberg Model "S" Air-Sea Dynamic Gravimeter, 1998; with permission

LaCoste/Romberg TAGS-6

(Turn-key Airborne Gravity System)

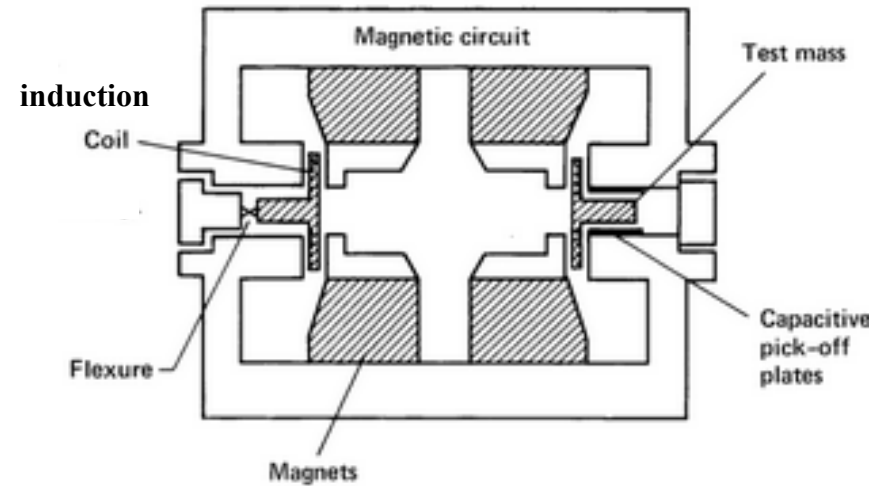


<http://www.microglacoste.com/tags-6.php>

SPECIFICATIONS		
COMPONENT	VARIABLE	SPECIFICATIONS
SENSOR	WORLDWIDE RANGE:	20,000 milliGal
	DYNAMIC RANGE:	±500,000 milliGal
	DRIFT:	3 milliGal per month or less
	TEMPERATURE SETPOINT:	45° to 65°C
STABILIZED PLATFORM	PLATFORM PITCH:	± 25 degrees
	PLATFORM ROLL:	± 35 degrees
	CONTROL:	
	Period	4 to 4.5 Minutes
Damping	0.707 of critical	
CONTROL SYSTEM	RECORDING RATE:	20 Hz
	SERIAL OUTPUT:	RS-232
	ADDITIONAL I/O:	Sensor Temperature
SYSTEM PERFORMANCE	DYNAMIC RANGE:	25,000,000
	STATIC REPEATABILITY:	0.02 milliGal in 2 min
	DYNAMIC REPEATABILITY:	0.75 milliGal in 2 min
MISCELLANEOUS	OPERATING TEMPERATURE:	5° to 50°C
	STORAGE TEMPERATURE:	-10° to 50°C
	POWER EQUIPMENTS:	75W @ 27°C Nominal 300W Peak
		80-265VAC, 47 – 63Hz
	DIMENSIONS:	58.4 x 53.3 x 55.9 cm (not including electronics)

BGM-3 Gravimeter

- Bell Aerospace (now Lockheed Martin)
 - Model XI **pendulous force-rebalance** accelerometer
 - **current needed to keep test mass in null position is proportional to acceleration**



(Seiff and Knight 1992); see also (Bell and Watts 1986)



Fugro → WHOI

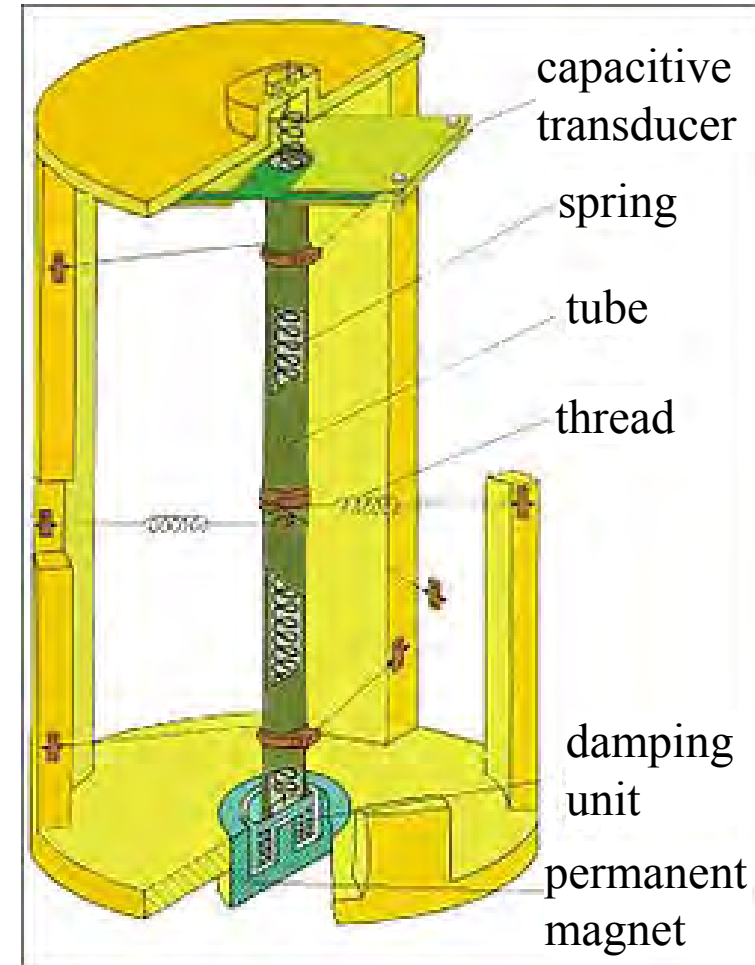
https://www.unols.org/sites/default/files/Gravimeter_Kinsey.pdf



Two BGM-3 gravimeters installed on the USCG ship *Healy*

Sea Gravimeter KSS31, Bodenseewerk Geosystem GmbH

- Gravity sensor based on Askania vertical-spring gravimeter
 - Federal Institute for Geosciences and Natural Resources (BGR)
 - **force rebalance** feedback system
 - highly damped output, ~ 3 minute average
 - sensor on a gyro-stabilized platform
 - also used for fixed wing and helicopter gravimetry



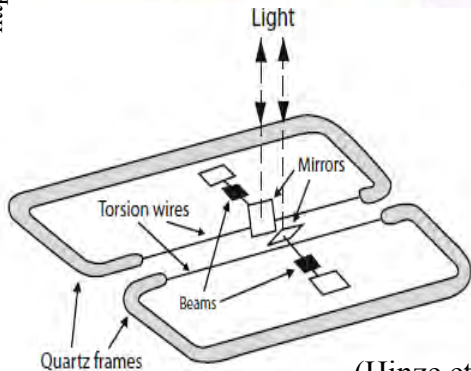
- http://www.bgr.bund.de/DE/Themen/MarineRohstoffforschung/Meeresforschung/Geraete/Gravimeter/gravimeter_inhalt.html
- http://www.bgr.bund.de/EN/Themen/GG_Geophysik/Aerogeophysik/Aerogravimetrie/aerogravimetrie_node_en.html

Chekan-A Gravimeter

- **Air-Sea gravimeter; CSRI* Elektropribor, St. Petersburg, Russia**
 - gravity sensed by deflection of pendulum hinged on quartz torsion wire in viscous fluid
 - pendulum deflection: 0.3–1.5 "/mGal; e.g., $\pm 1^\circ \rightarrow \pm 10$ Gal total range (0.36 "/mGal)
 - cross-coupling effect minimized by double-beam reverted pendulums
 - evolutionary modifications: Chekan AM, "Shelf" (Krasnov et al. 2014)

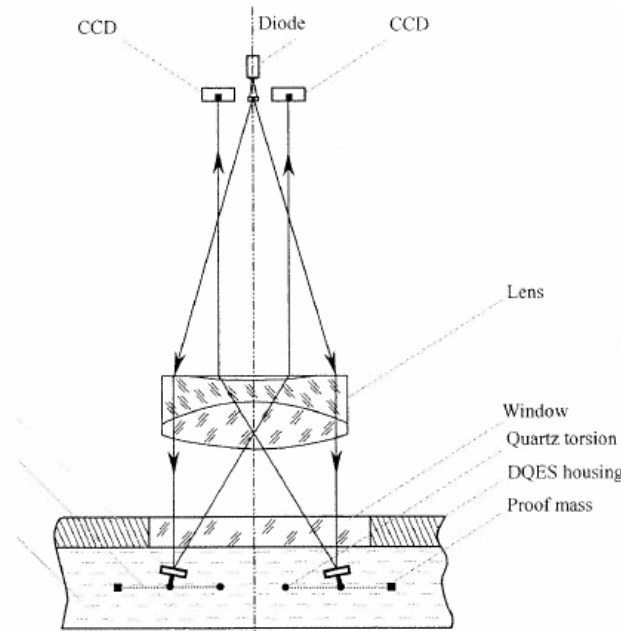
*Central Scientific & Research Institute

http://www.gravionic.com/gravimetry.html



(Hinze et al. 2013)

Pendulum mirror
 Pendulum lever
 Damping liquid



(Krasnov et al. 2008)

Installation in test aircraft



(Stelkens-Kobsch 2005)

Airborne Gravimetry on Airship Platform



(<http://rosaerosystems.com/airships/obj17>)

Relative gravimeter Chekan in a cabin of AU



– **test flight January 2014**

– **reported in IAG Commission 2 Travaux 2015**

Airborne Gravimeter GT-2A

- Gravimeter system designed by *Gravimetric Technologies (Russia)*
 - vertical accelerometer of axial design with a test mass on spring suspension
 - photoelectric position pickup
 - moving-coil force feedback transducer
 - three-axis gyro-stabilized platform
 - large dynamic range



(Gabell et al. 2004)



Measurement range	9.75 to 9.85 m/sec ²
Dynamic range	> +/- 1,000 Gals
Drift per day (corrected)	< 0.1 mGals
RMS error in gravity anomaly estimation (static mode up to 12 hours on bench)	
RMS error	0.6 mGals (+/- 1 LSD*)
Attitude limits	
roll	+/- 45°
pitch	+/- 45°
Operating temp	+5°C to +50°C
Power	
operating	150 W at 27Vdc
standby	50 W at 27Vdc
Weight (with base)	153.5 kg
Dimensions console	400 x 400 x 600 mm
Dimensions base	600 x 300 mm
Service life	30,000 hours
Error in gravity anomaly estimation (RMS)	
0.01 Hz cut-off	0.6 mGals (+/- 1 LSD*)
*Least Significant Digit Specifications subject to change	

http://eongeosciences.com/wp-content/uploads/2015/01/GT_2A.pdf
(Canadian Micro Gravity)

Sander Geophysics Ltd. AirGrav System

- “Purpose-built” airborne gravimeter designed for airborne environment, not modified sea gravimeter
 - Honeywell inertial navigation grade **accelerometers** (Anecchione et al. 2006, Sinkiewicz et al. 1997)
 - Schuler-tuned (three axis) inertially stabilized platform
 - three-accelerometer system; vertical accelerometer used as **gravimeter**
 - advertize gravimetry on **topography-draped** profiles
 - demonstrated success in **vector** gravimetry
 - comparison tests over Canadian Rockies between AirGrav and GT-1A (Studinger et al. 2008)



<http://www.sgl.com/news/Sander%20Geophysics%20-%20Antarctica.pdf>